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Directly detecting the Milky Way halo

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Kings College London Dec 2016

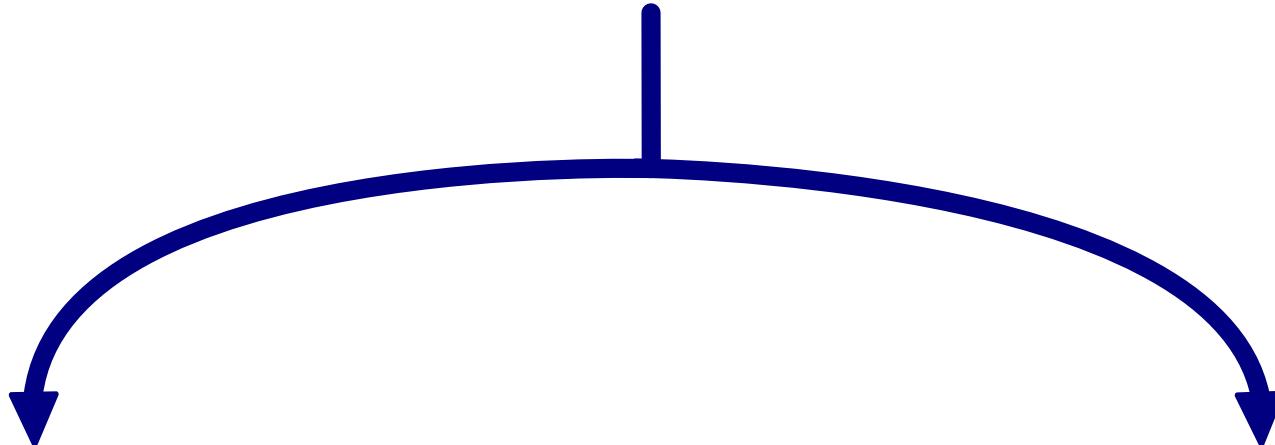
Outline

Astrophysical uncertainties in dark matter detection

➤ C. A. J. O'Hare [1604.03858]



Observing the local Milky Way dark matter distribution



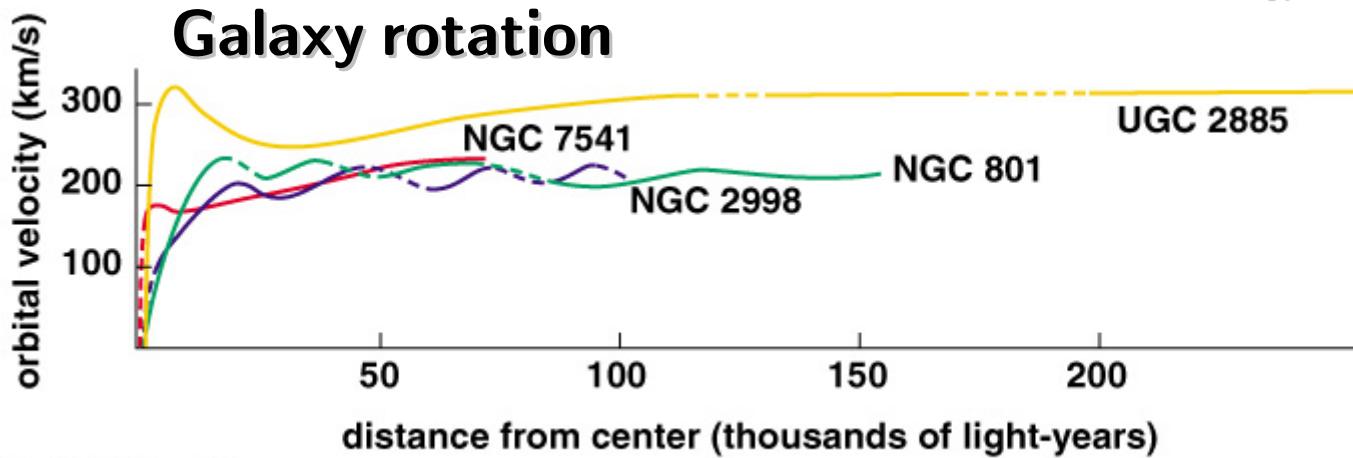
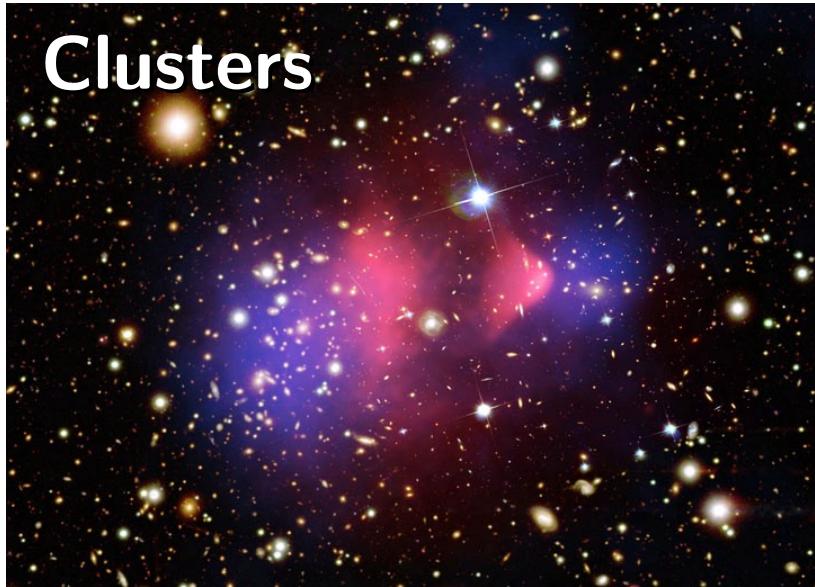
WIMP directional detectors

- C. A. J. O'Hare & A. M. Green [1410.2749]
- B. J. Kavanagh & C. A. J. O'Hare [1609.08630]

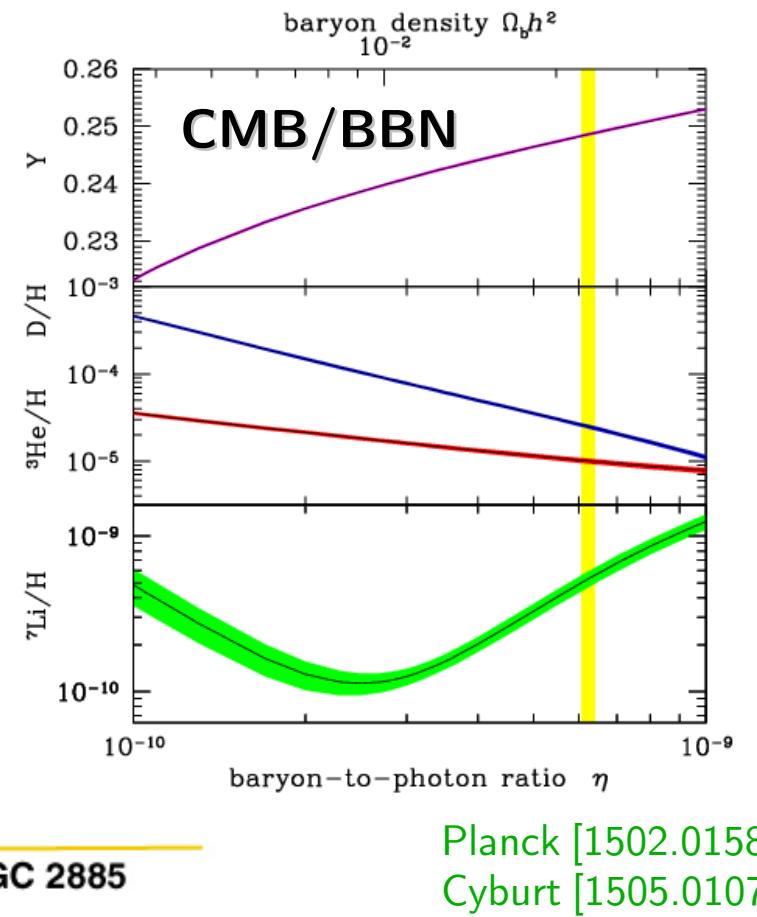
Axion haloscopes

- C. A. J. O'Hare & A. M. Green [in prep.]

We haven't enough baryons...



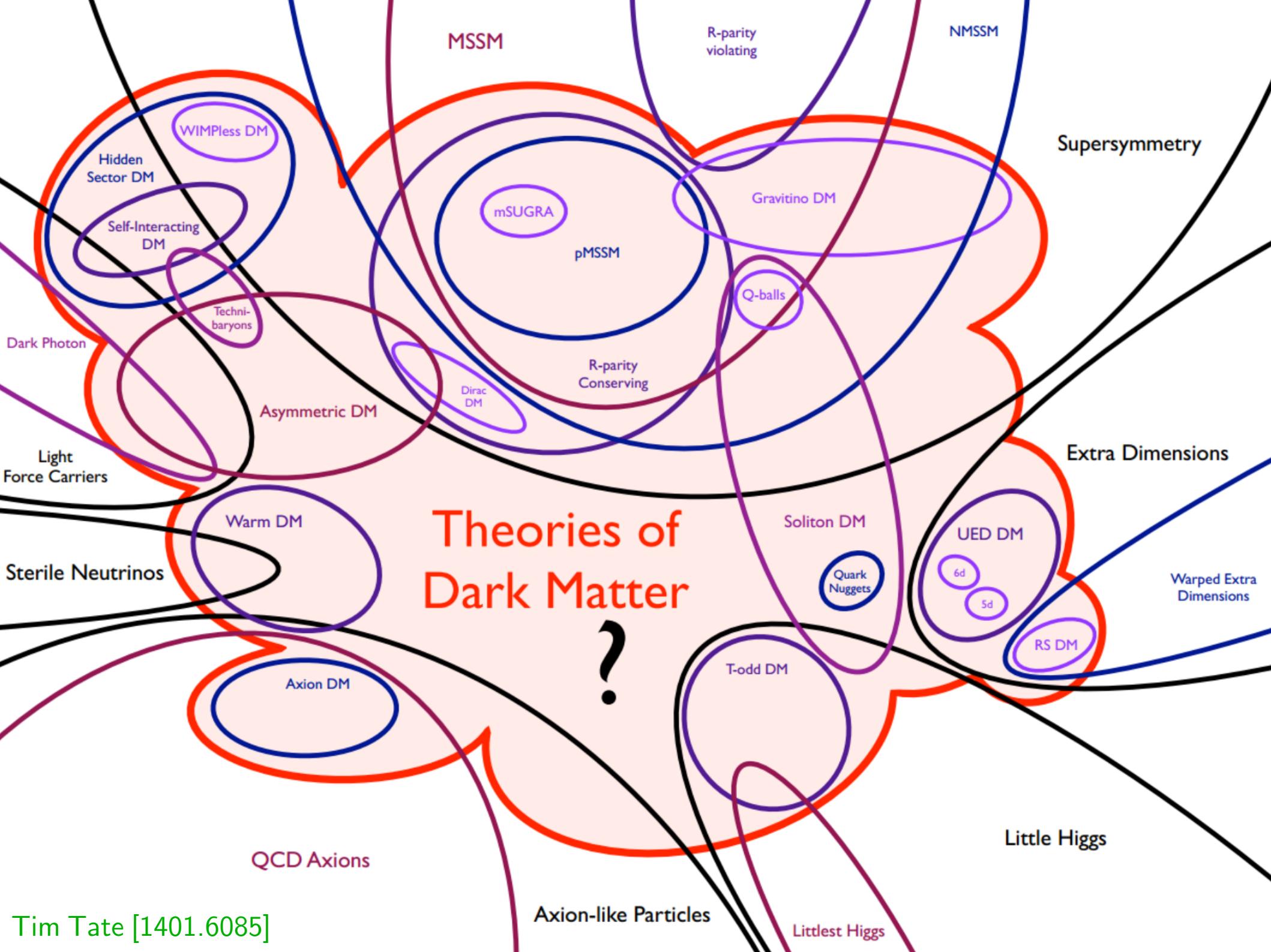
(c) Addison Wellsley (2009)



...dark matter is there and it's probably a particle

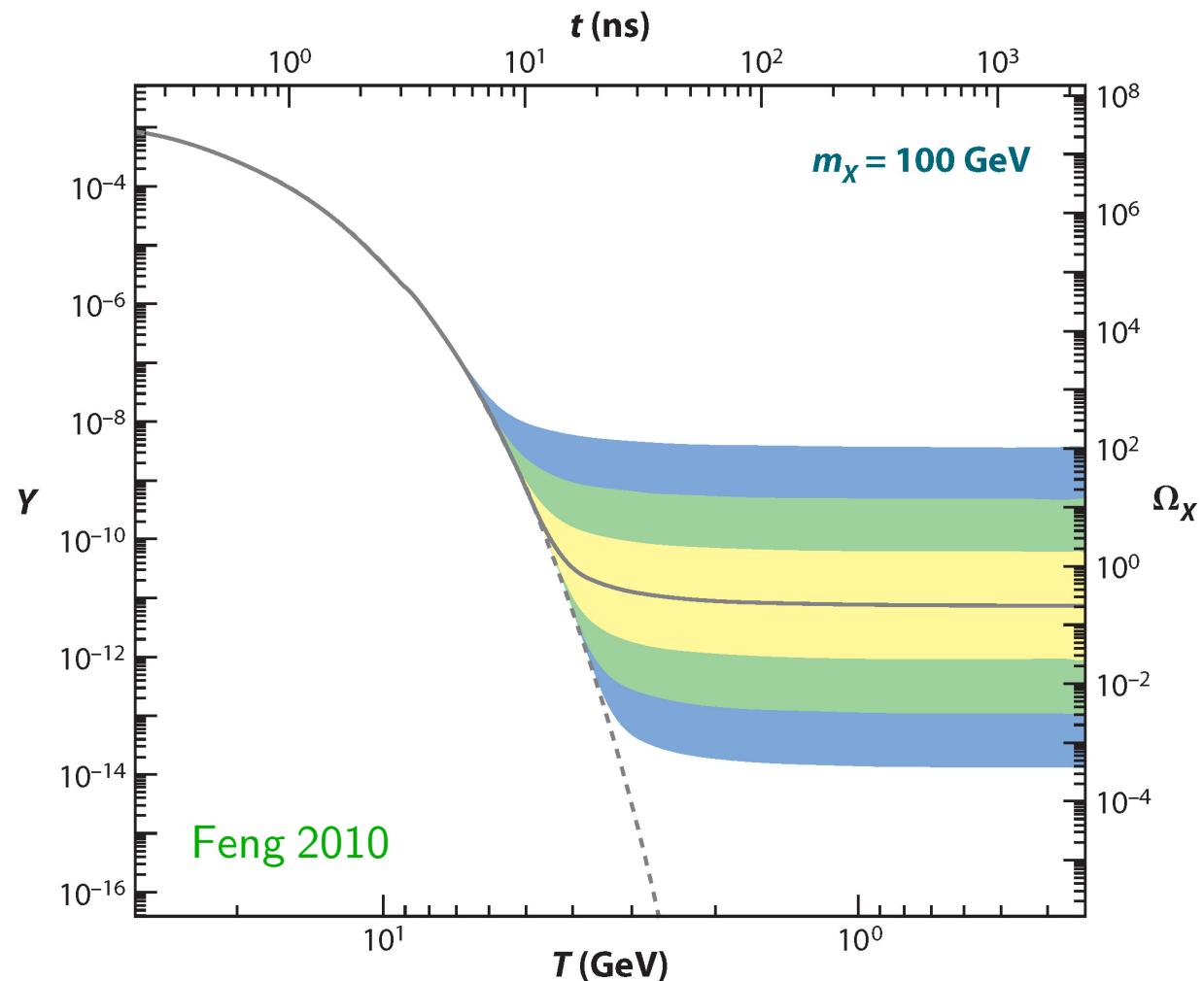
Theories of Dark Matter

?



WIMPs

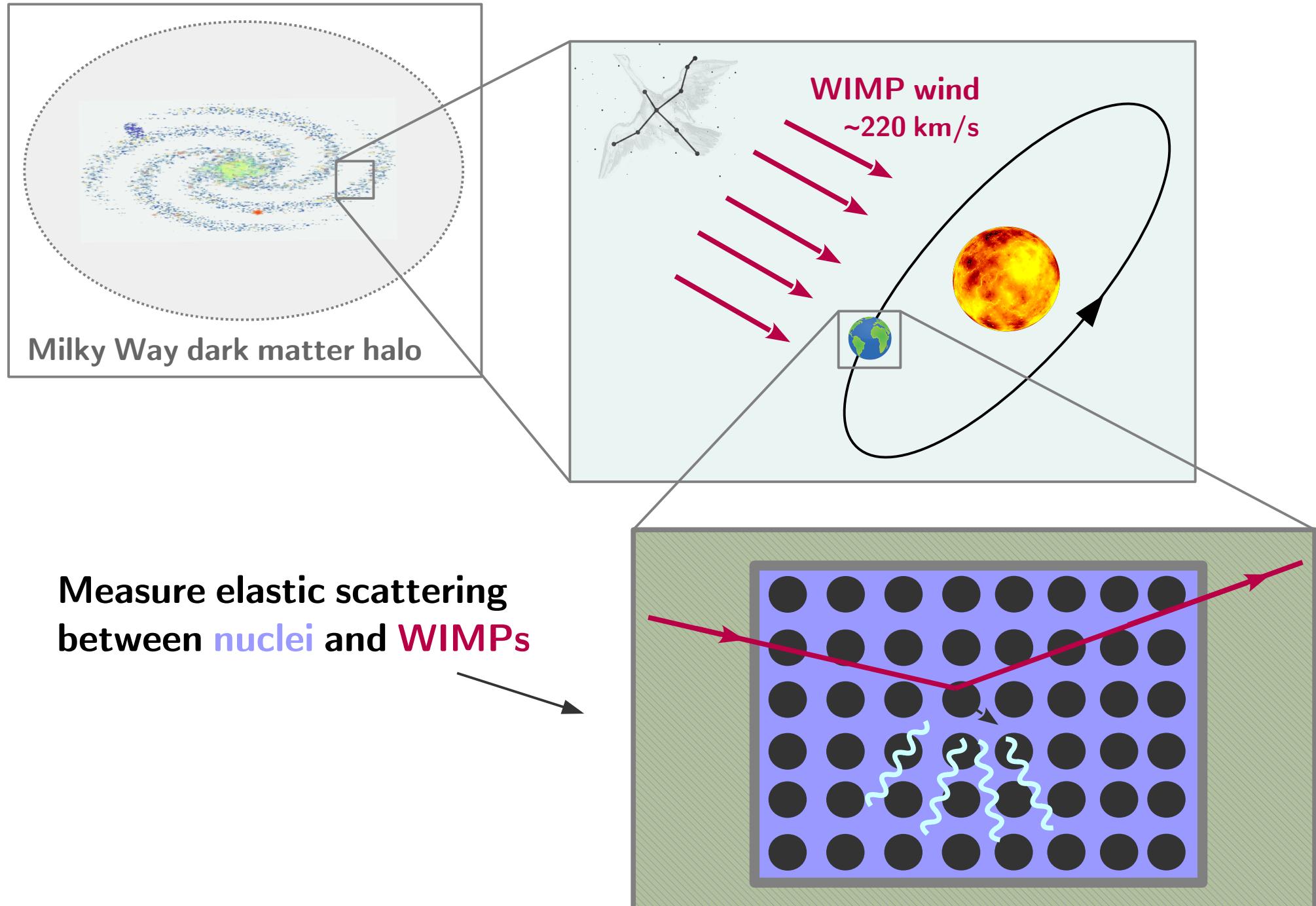
Assuming new particle is initially in thermal equilibrium
→ freeze-out with relic density proportional to 1/annihilation cross section



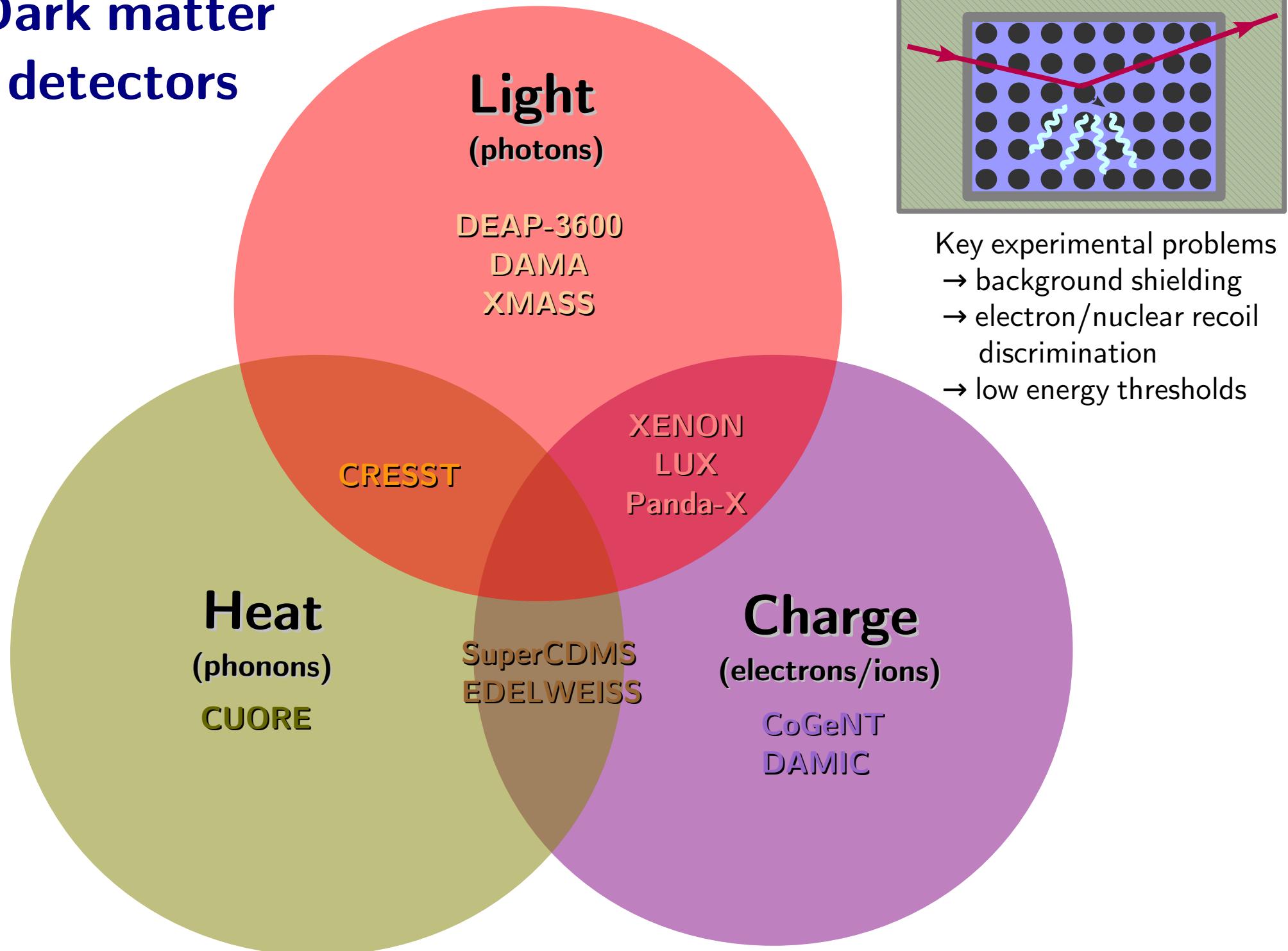
$$\Omega_{\text{dm}} h^2 \simeq 0.1196 \implies \langle \sigma v \rangle \simeq 3 \times 10^{-26} \text{ cm}^3 \text{ s}^{-1}$$

- Observed relic density would imply weak-scale annihilations
- WIMPs show up in BSM physics

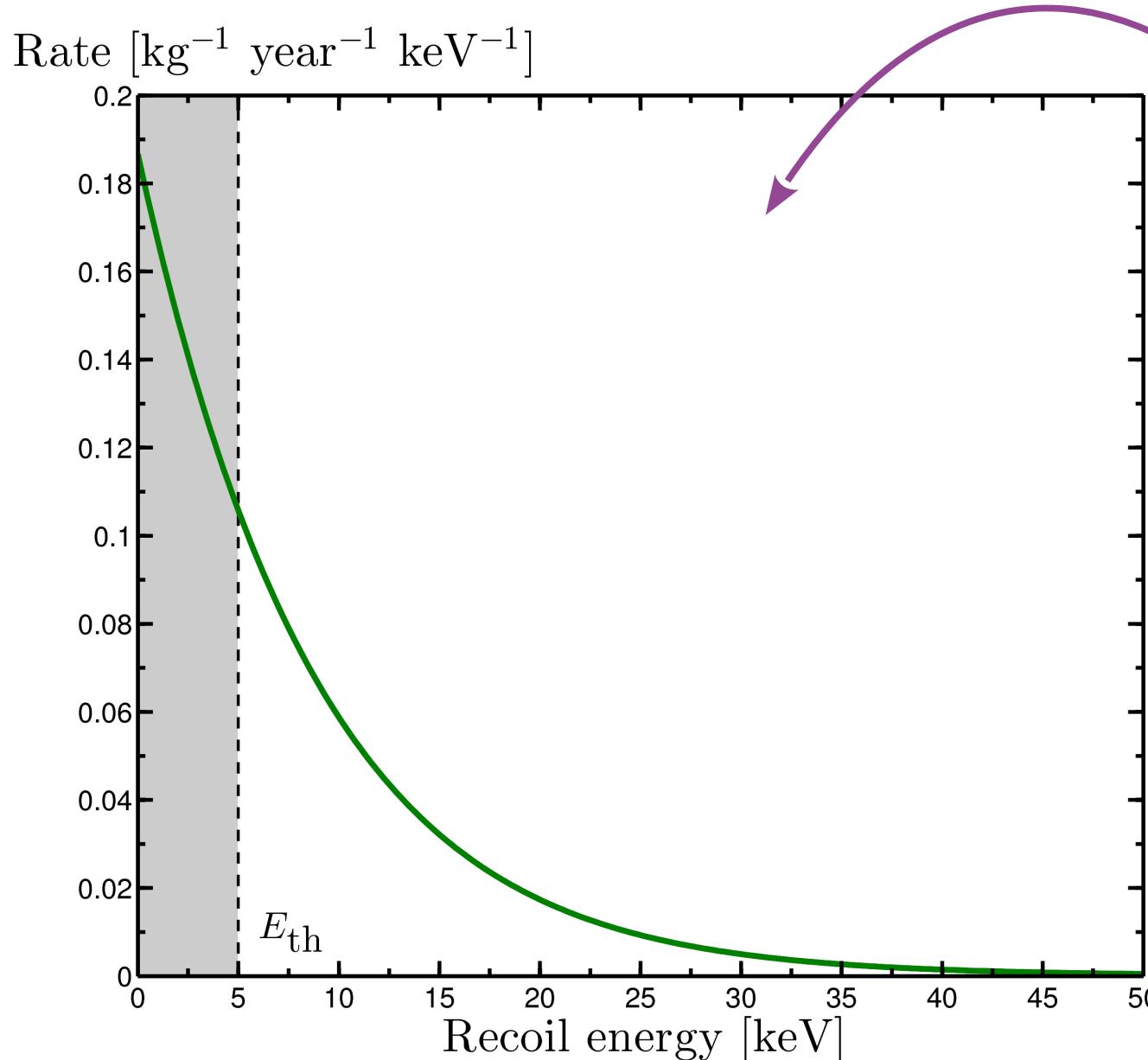
Direct detection



Dark matter detectors

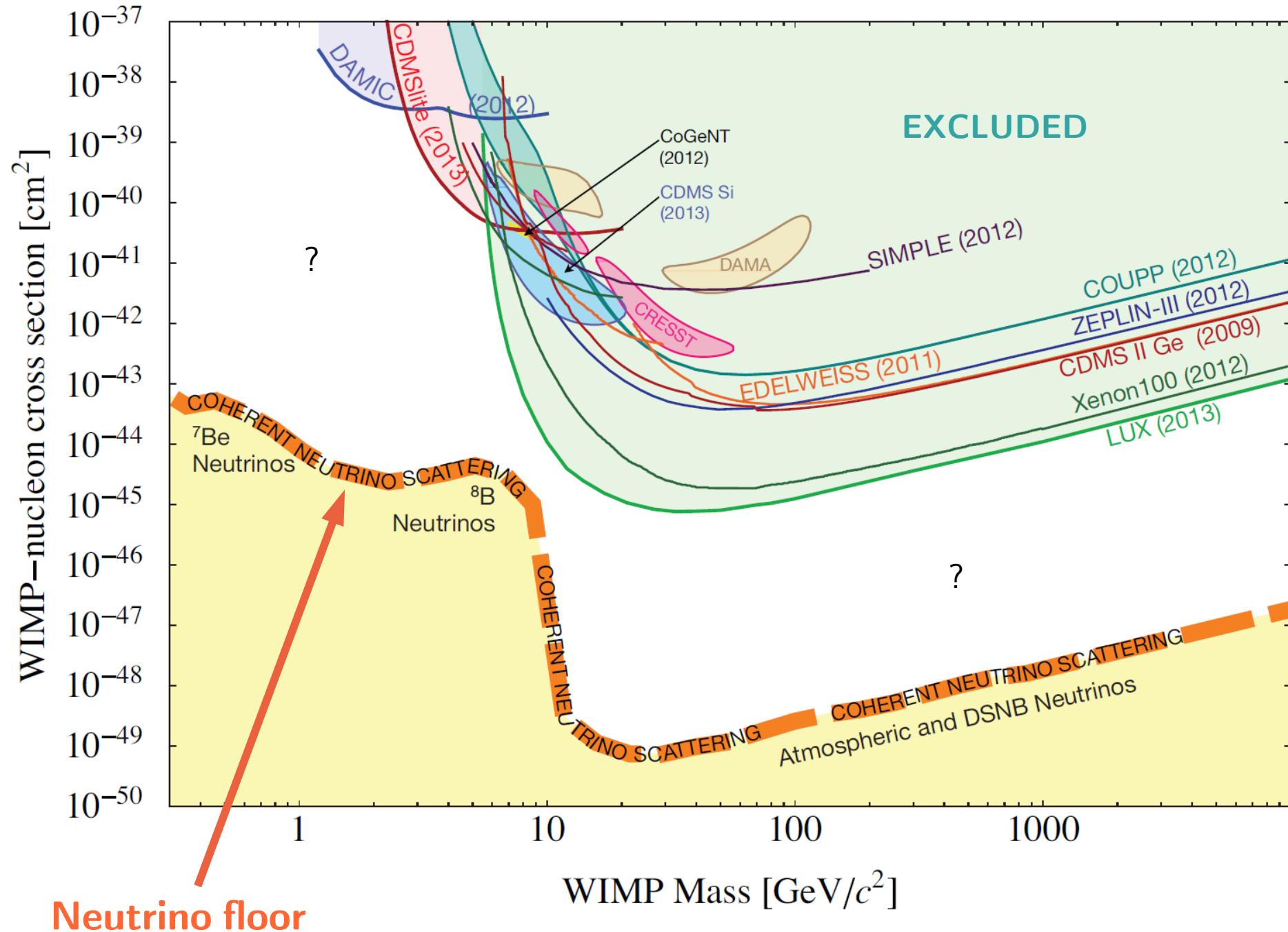


Direct dark matter detection



Encodes:

- WIMP mass & cross section
- Nuclear response
- Detector effects
- **Astrophysical information...**



WIMP-nucleus scattering

Double differential event rate (recoil energy E_r and direction \mathbf{q})

$$\frac{dR}{dE_r d\Omega_q} = \frac{\sigma_p \mathcal{C}_N}{4\pi \mu_{\chi p}^2 m_\chi} \times F^2(E_r) \times \rho_0 \hat{f}(v_{\min}, \hat{\mathbf{q}})$$

Particle physics

- WIMP-nucleus interaction
 - spin-independent: $\mathcal{C}_N^{\text{SI}} \propto A^2$
 - spin-dependent: $\mathcal{C}_N^{\text{SD}} \propto (J+1)/J$

Nuclear physics

- Form factor $F^2(E_r)$

Astrophysics

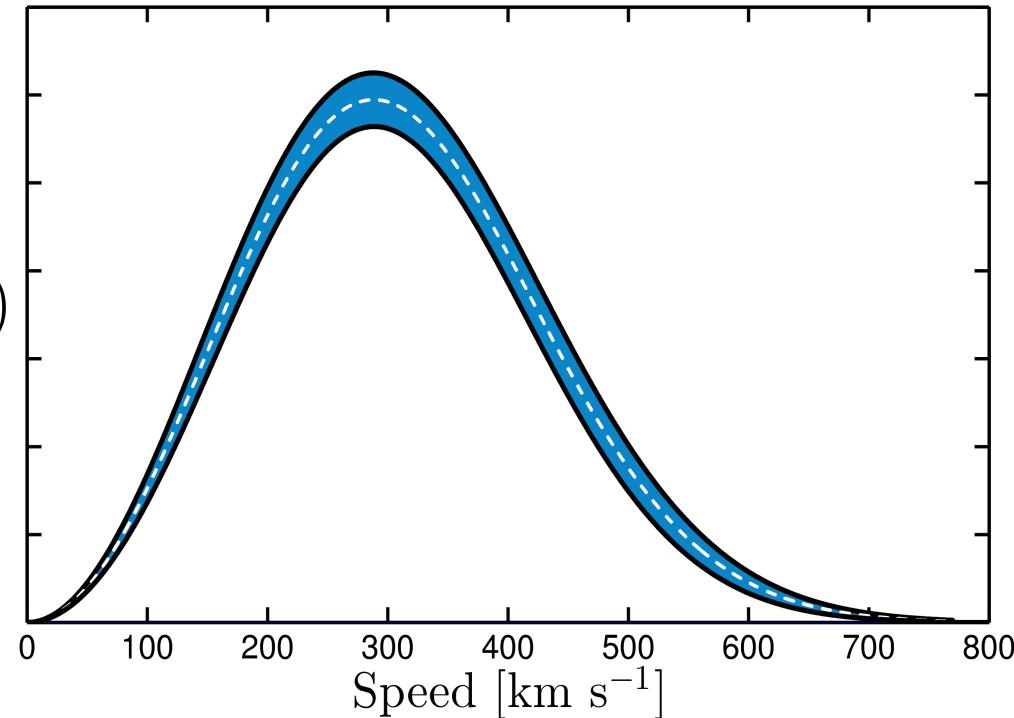
- Local DM density ρ_0
- Radon transform of velocity distribution:

$$\hat{f}(v_{\min}, \hat{\mathbf{q}}) = \int f_{\text{lab}}(\mathbf{v}, t) \delta(\mathbf{v} \cdot \hat{\mathbf{q}} - \hat{v}_{\min}) d^3 \mathbf{v}$$

Astrophysical uncertainties

- Experimental analyses typically assume **Standard Halo Model (SHM)**
- Smooth isothermal sphere (cow)
→ Maxwell-Boltzmann velocity distribution

$$f(\mathbf{v}) = \begin{cases} \frac{1}{N} e^{-v^2/v_0^2} & |\mathbf{v}| < v_{\text{esc}} \\ 0 & |\mathbf{v}| \geq v_{\text{esc}} \end{cases}$$



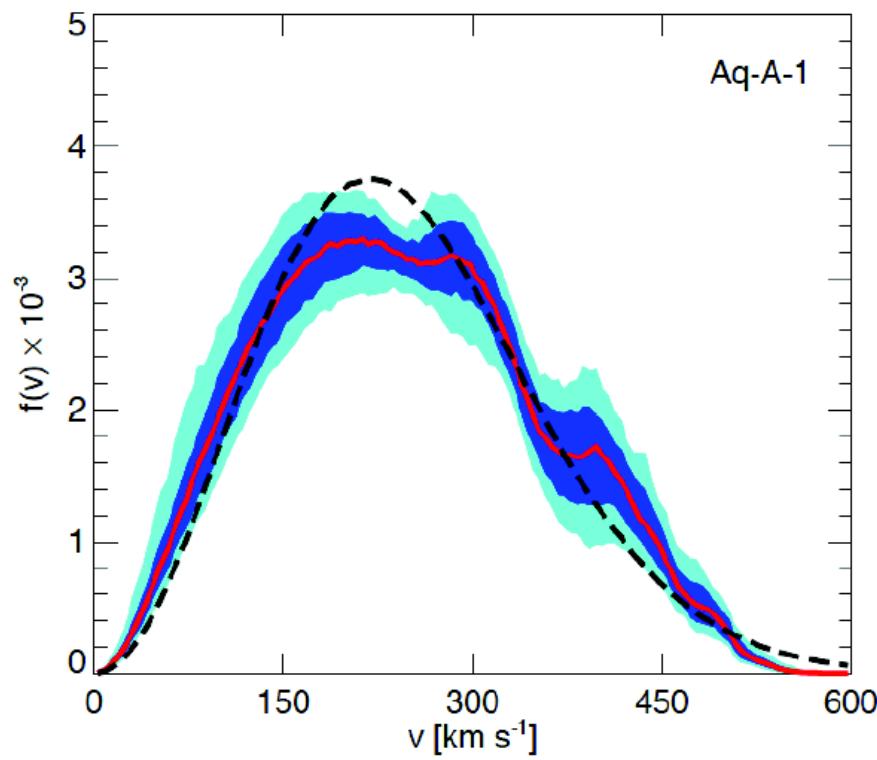
Free parameters:

- Rotation speed, $v_0 = 180 - 240 \text{ km s}^{-1}$ [Lavalle & Magni \[1411.1325\]](#)
- Escape velocity, $v_{\text{esc}} = 533 \pm 50 \text{ km s}^{-1}$ [Piffl et al \[1309.4293\]](#)
- Local density, $\rho_0 = 0.2 - 0.8 \text{ GeV cm}^{-3}$ [Read et al \[1404.1938\]](#) ...

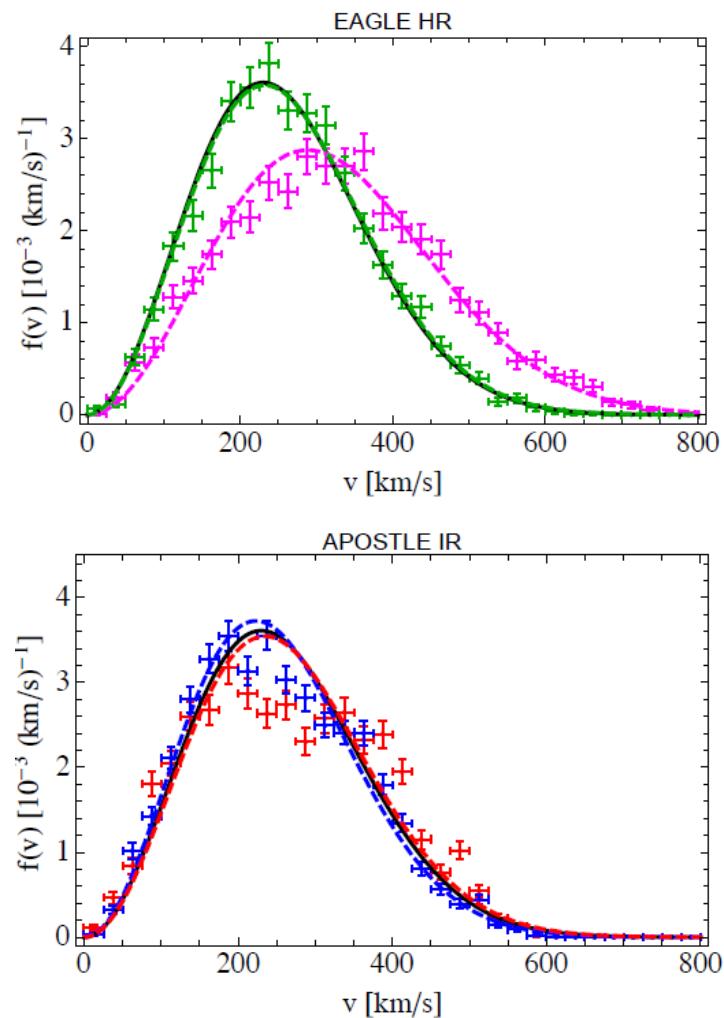
Halo simulations

- N-body/hydrodynamic simulations persistently exhibit non-Maxwellian structure

*Although adding baryons can improve Maxwellian fit



Vogelsberger [0812.0362]

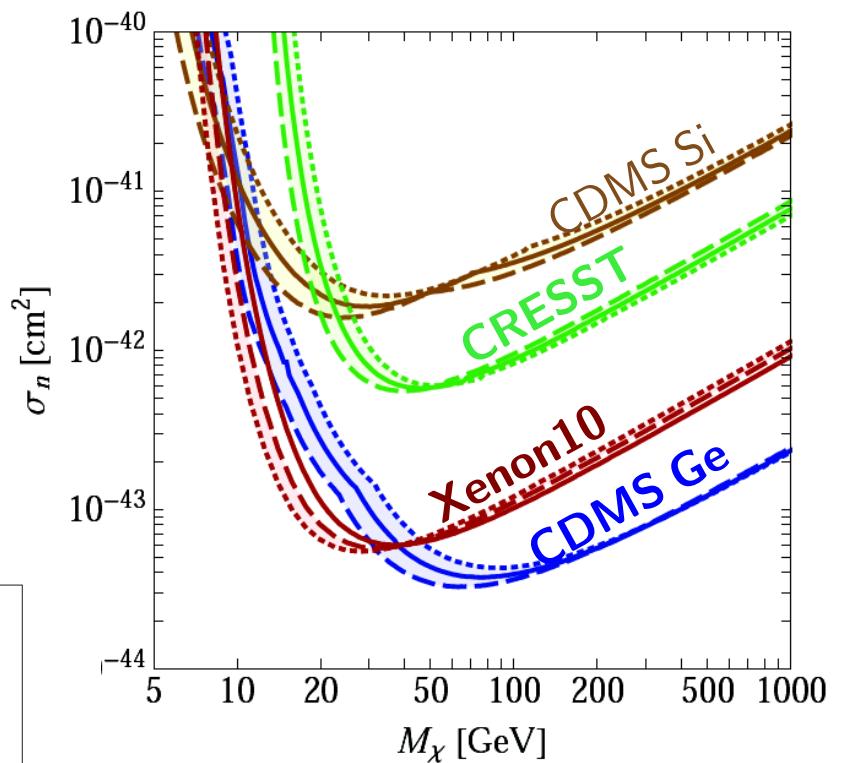
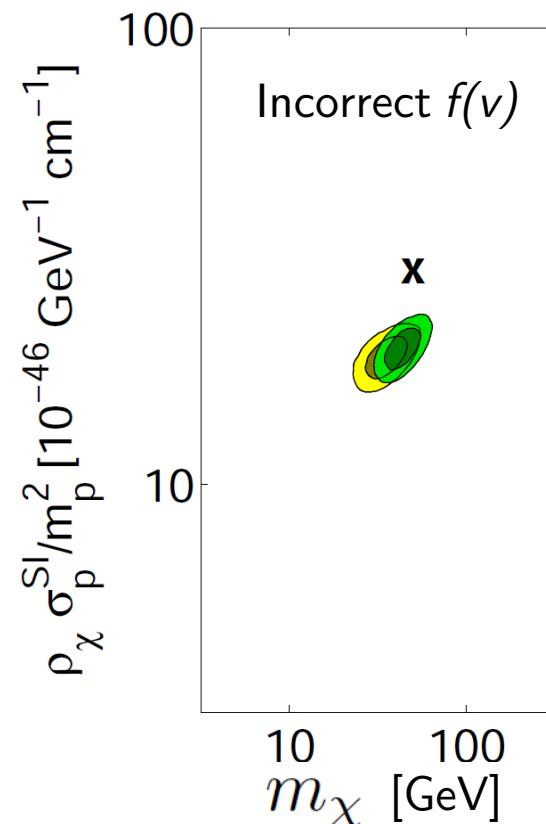
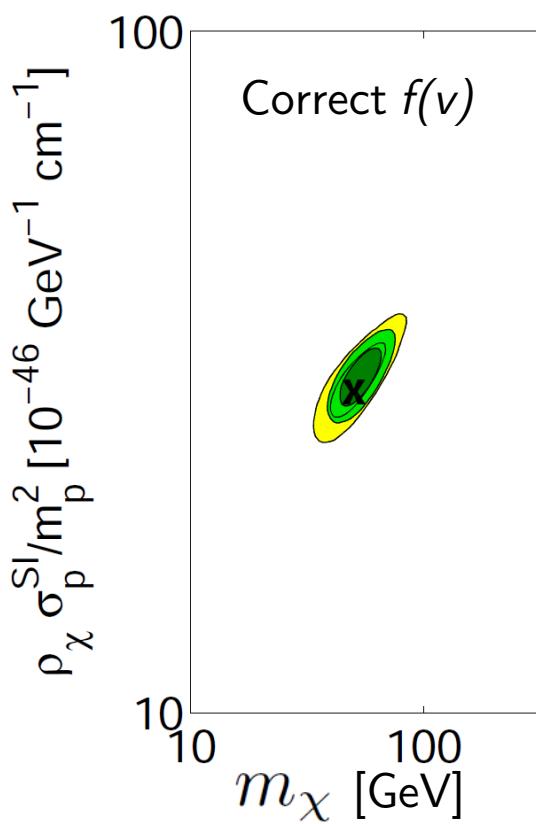


Calore [1509.02164]

Bozorgnia [1601.04707]

Effect of astrophysical uncertainties

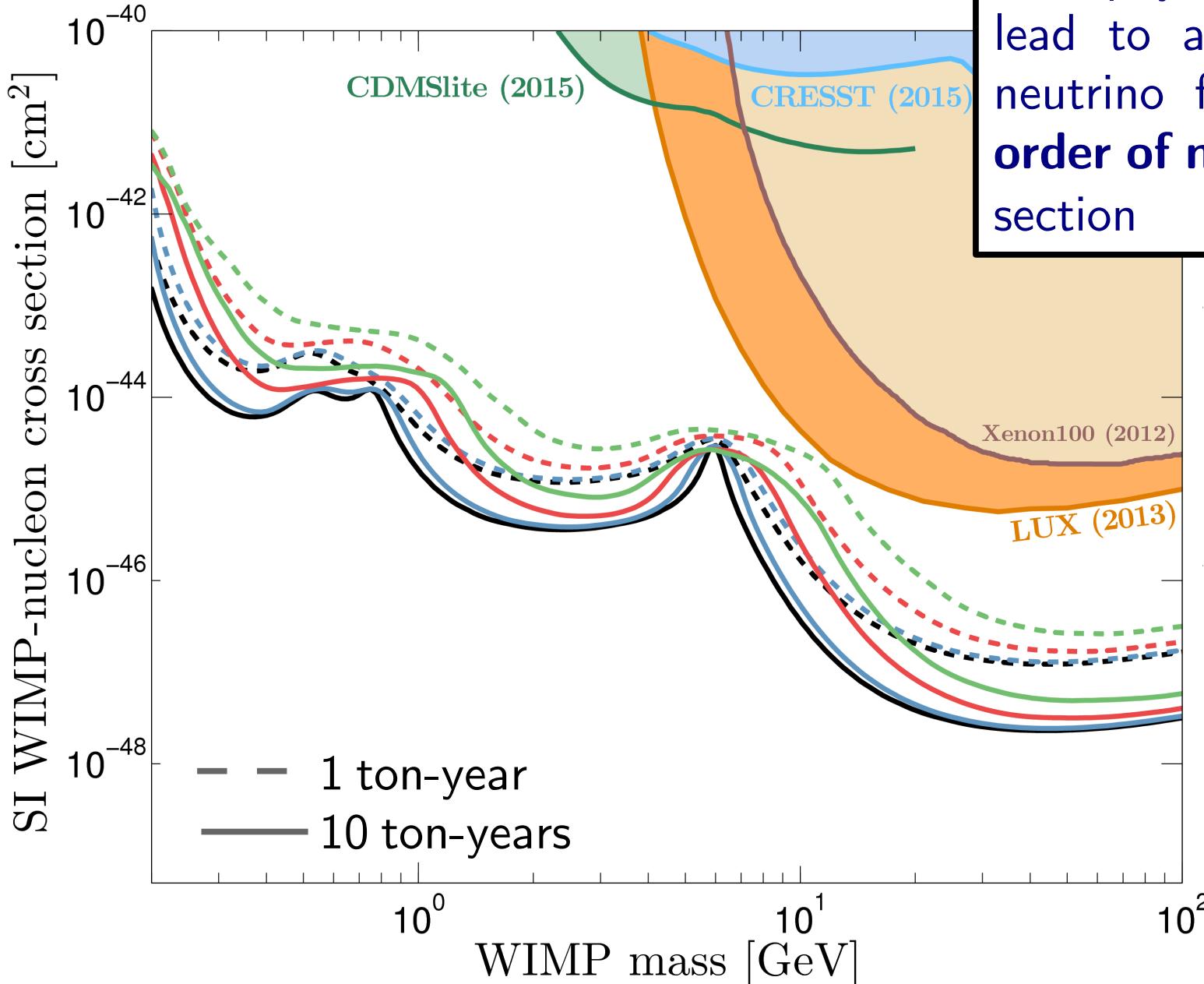
Uncertainty in exclusion
limits →
e.g. McCabe [1005.0579]



← Biased parameter
estimation
e.g. Peter [1103.5145]

The neutrino floor

CAJ O'Hare [1604.03858]



Astrophysical uncertainties lead to an increase in the neutrino floor of up to an **order of magnitude** in cross section

Low: $\sigma_{\rho_0} = 0.01 \text{ GeV cm}^{-3}$
 $\sigma_{v_0} = 10 \text{ km s}^{-1}$
 $\sigma_{v_{\text{esc}}} = 10 \text{ km s}^{-1}$

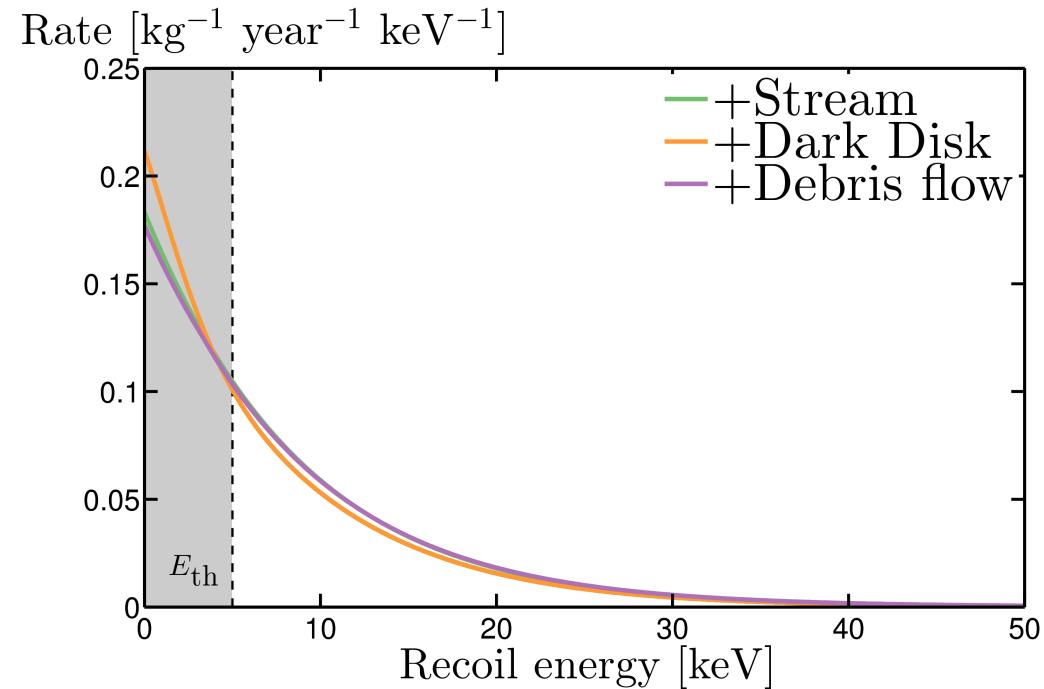
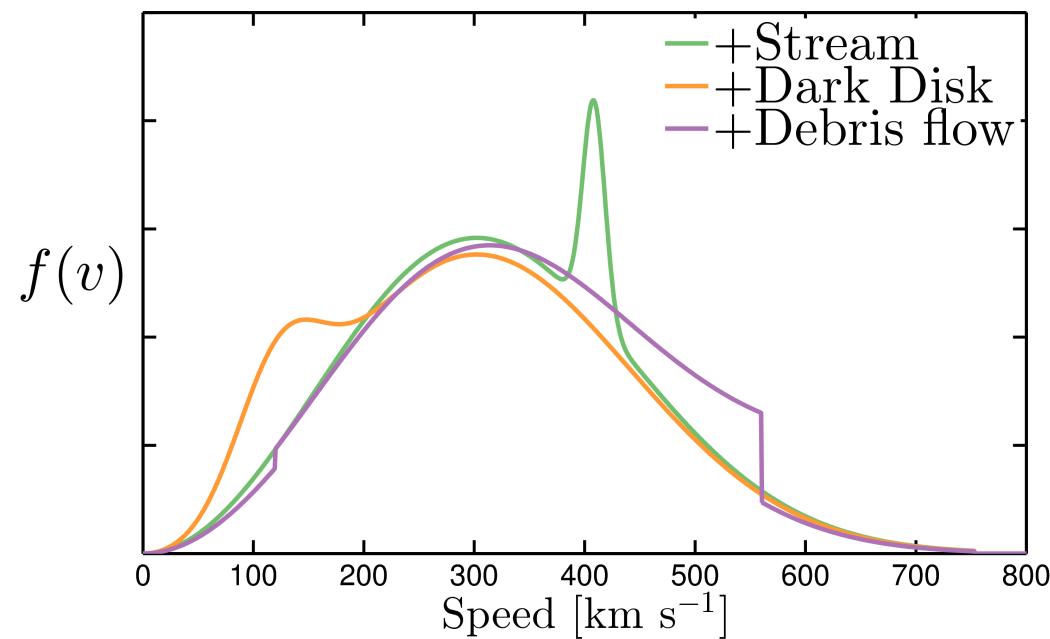
Med: $\sigma_{\rho_0} = 0.05 \text{ GeV cm}^{-3}$
 $\sigma_{v_0} = 40 \text{ km s}^{-1}$
 $\sigma_{v_{\text{esc}}} = 40 \text{ km s}^{-1}$

High: $\sigma_{\rho_0} = 0.1 \text{ GeV cm}^{-3}$
 $\sigma_{v_0} = 60 \text{ km s}^{-1}$
 $\sigma_{v_{\text{esc}}} = 50 \text{ km s}^{-1}$

— No uncertainties

Dark matter substructure

Observations/simulations suggest possible substructure in local MW

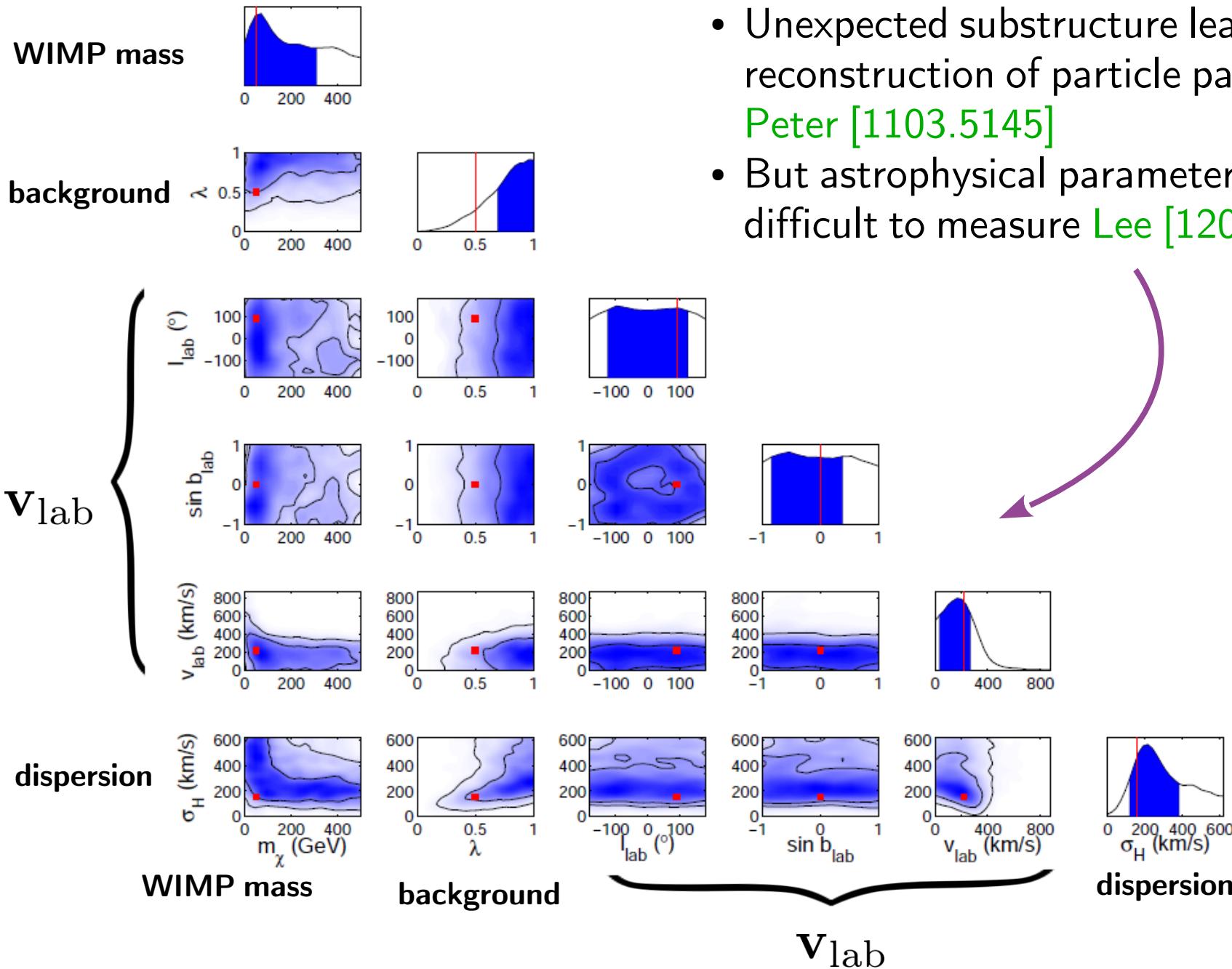


Speed distribution \longrightarrow Scattering rate

Possible substructures

- Tidal streams Purcell *et al.* [1203.6617]
- Dark disk Pillepich *et al.* [1308.1703], Schaller *et al.* [1605.02770]
- Debris flows Kuhlen *et al.* [1202.0007]

Measuring astrophysical parameters



Astrophysical uncertainties cause problems for direct detection:

- Uncertainty in experimental limits
- Neutrino floor higher
- Degeneracy with particle physics parameters
- Possible presence of substructure

**Solution: go and measure the local
Milky Way halo directly**

Astrophysical uncertainties cause problems for direct detection:

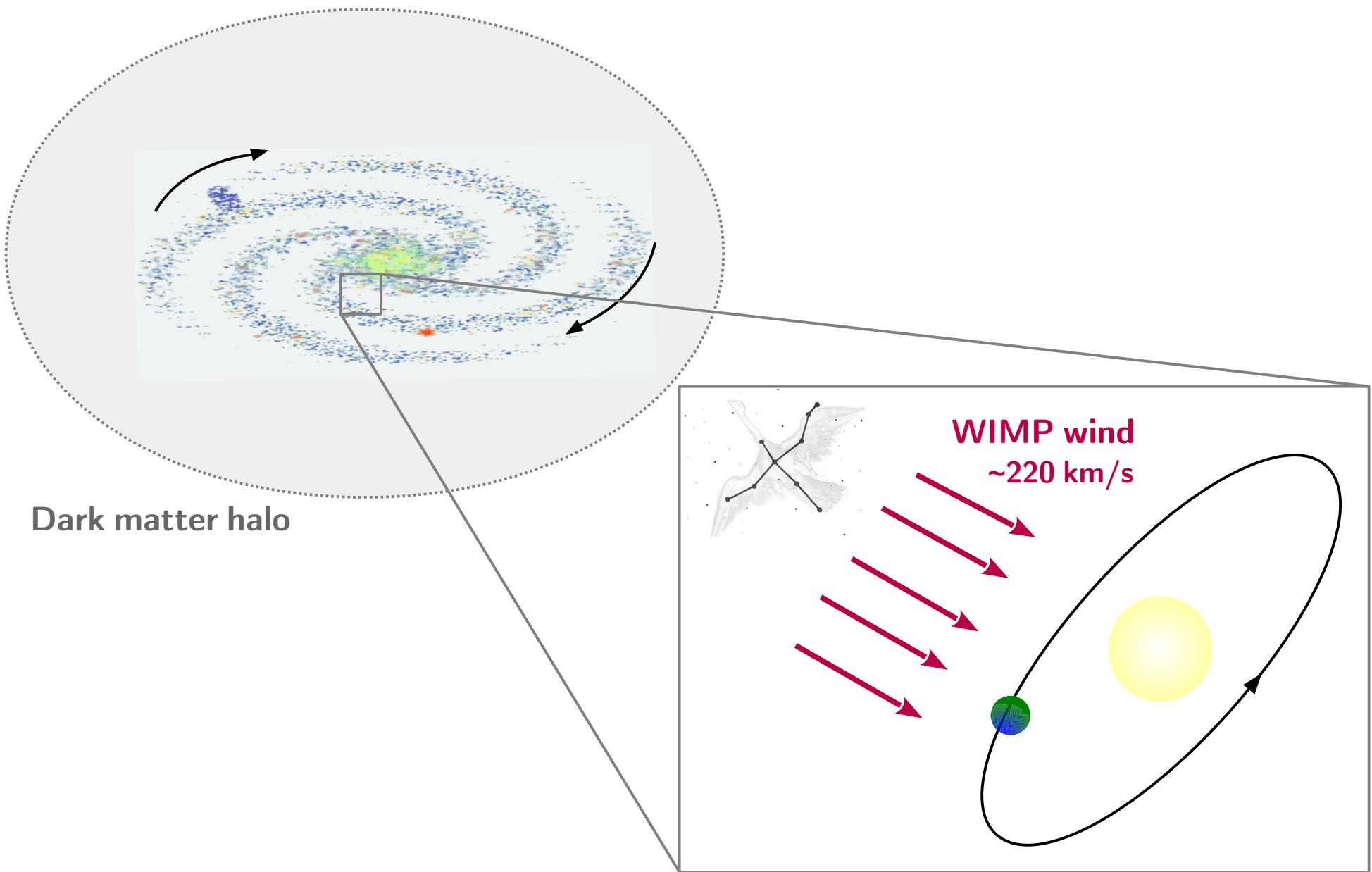
- Uncertainty in experimental limits
- Neutrino floor higher
- Degeneracy with particle physics parameters
- Possible presence of substructure

Solution: go and measure the local Milky Way halo directly

Bonus: find out about the formation history of the Milky Way...

Directional detection

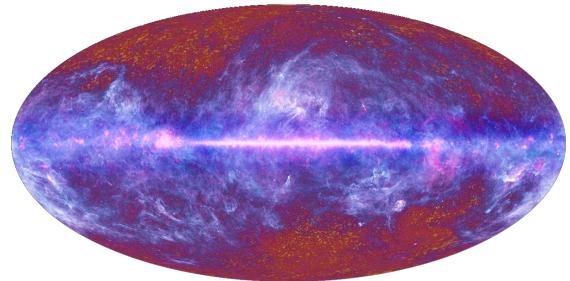
Directional detection



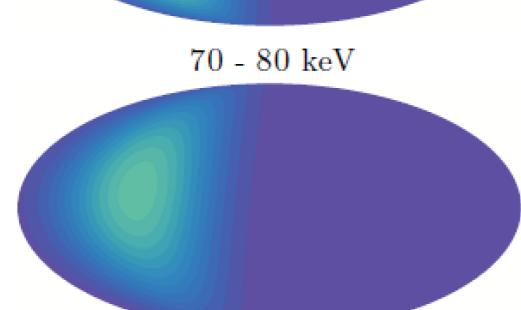
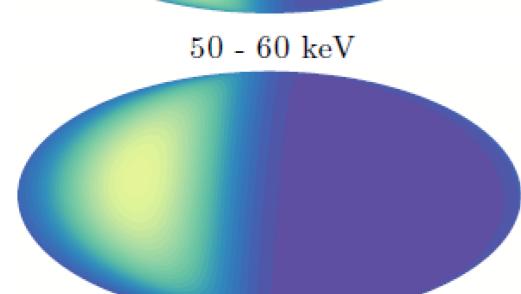
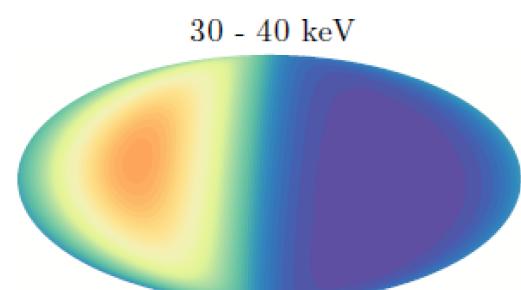
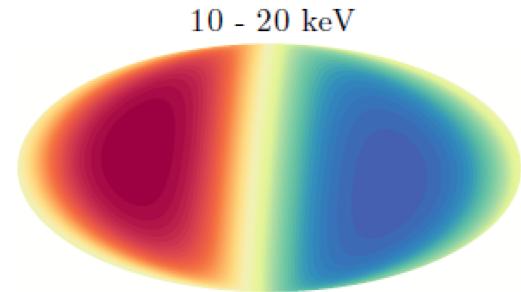
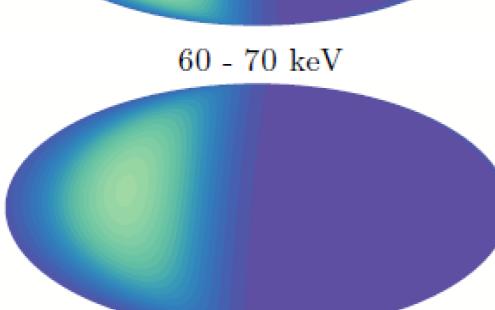
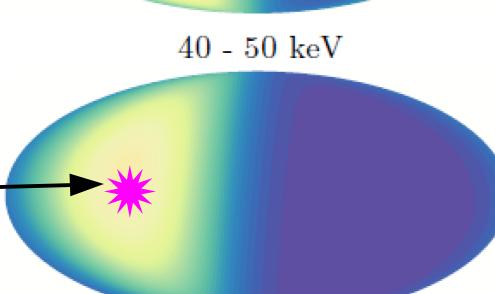
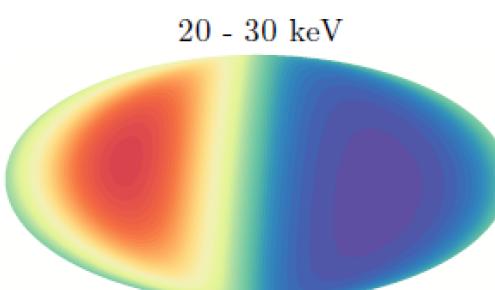
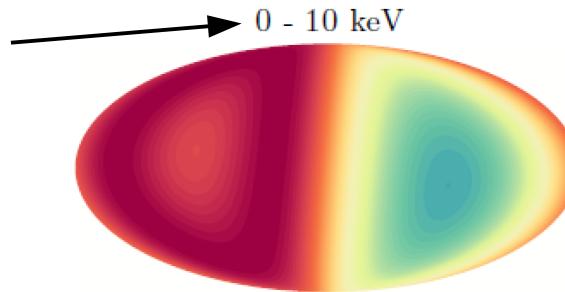
Energy+Direction dependence

100 GeV WIMP

^{19}F recoil energy range
Mollweide projection of recoil skymap
c.f.



Peak direction toward $-\mathbf{v}_{\text{lab}}$

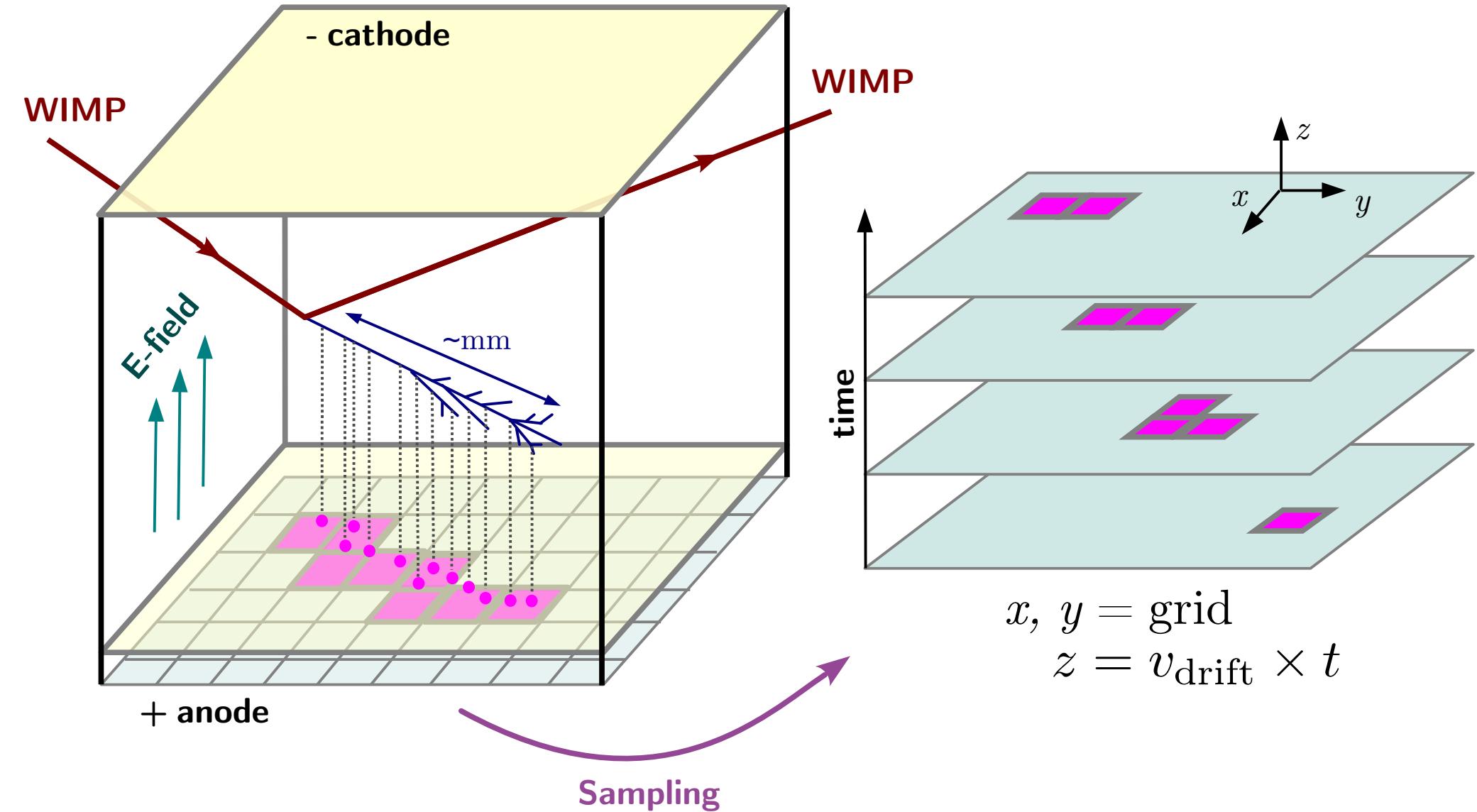


Secondary features

- Ring at low energies [Bozorgnia \[1111.6361\]](#)
- Aberration over time [Bozorgnia \[1205.2333\]](#)

Directional detection

- Low pressure gas time projection chamber (TPC):



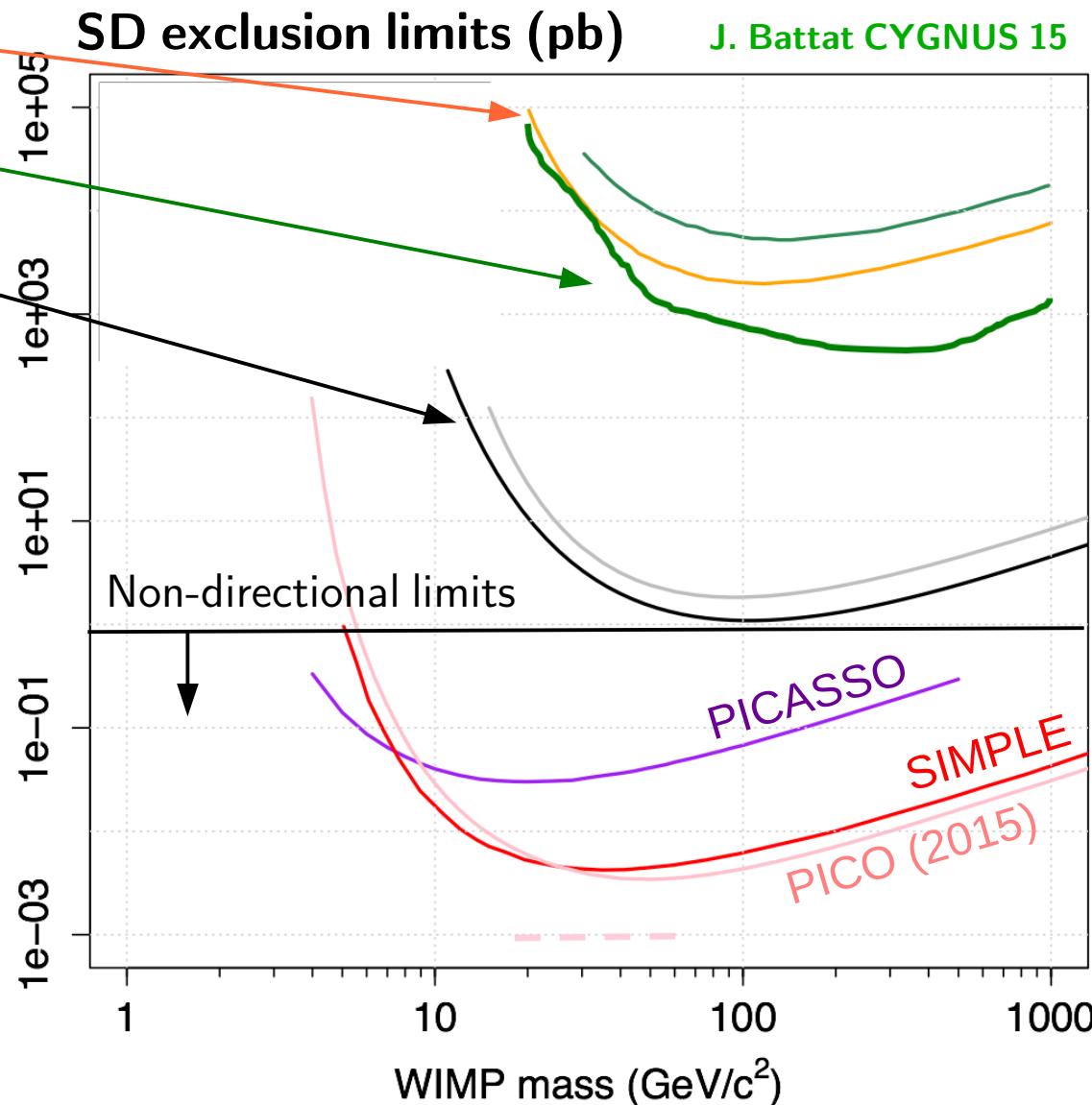
Directional detectors

Best at the moment: low pressure gas TPCs

- DM-TPC (USA)
- NEWAGE (Japan)
- DRIFT (UK)
- MIMAC (France)

Many other ideas*...

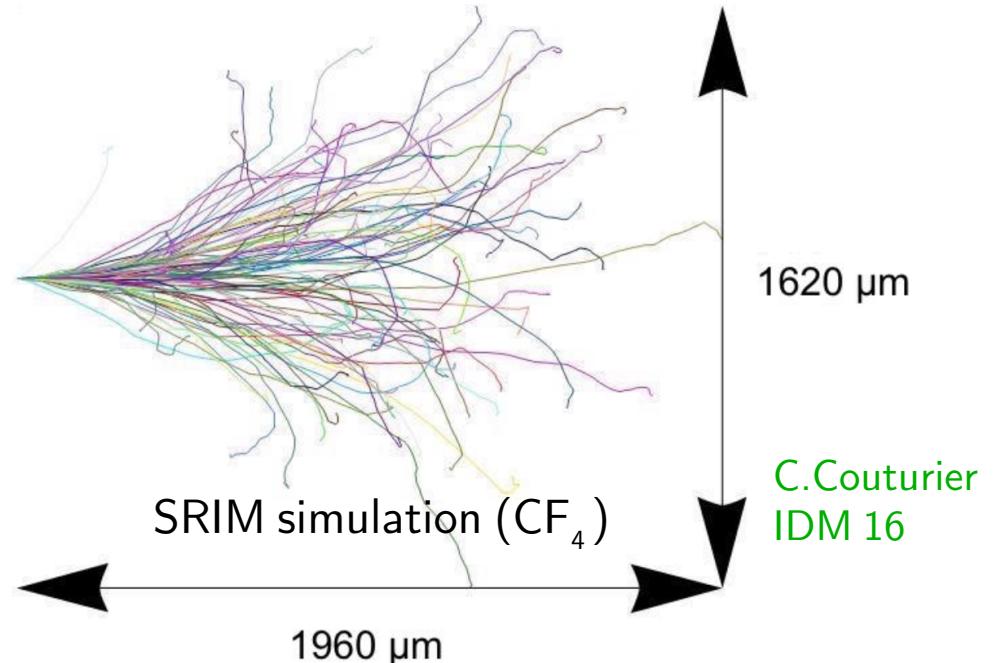
- Emulsion plates
- Crystal scintillators
- LXe/Ar Columnar recombination
- DNA (perhaps...)



Directional detection

Disadvantages (TPCs):

- Inherently low mass $O(0.1 \text{ kg})$
- Angular resolution $O(10) \text{ deg.}$
- Sense recognition ($+q$ or $-q$)
- 1-d/2-d/3-d readout

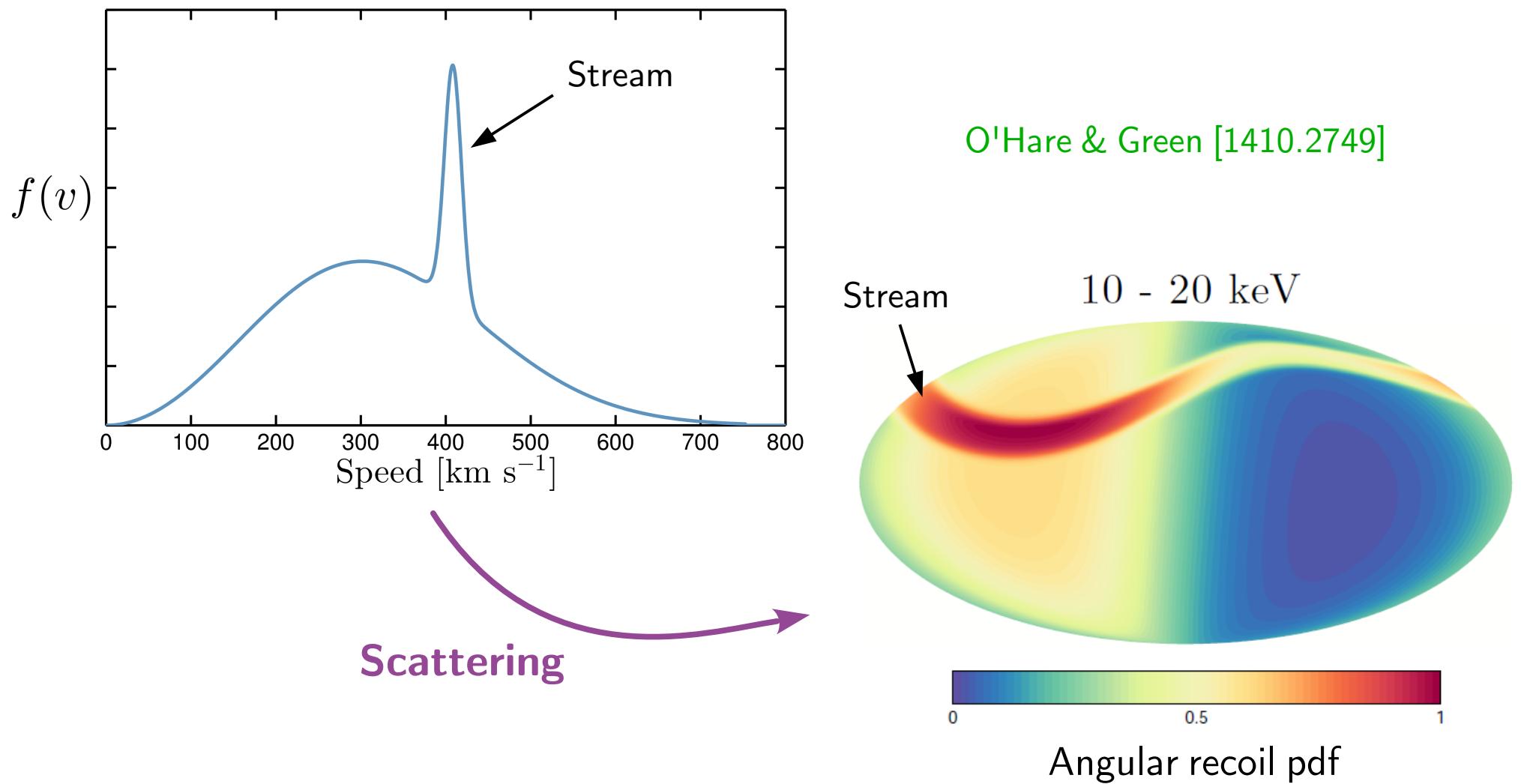


Advantages:

- Excellent electron-nuclear recoil discrimination Billard *et al* [1205.0973]
- Reject isotropy with $O(10)$ events Morgan *et al* [astro-ph/0408047]
- Discover DM with $O(30)$ events Green & Morgan [1002.2717]
- No neutrino floor Grothaus *et al* [1406.5047]
- Access *velocity* distribution Kavanagh & O'Hare [1609.08630]
- Probe DM substructure O'Hare & Green [1410.2749]

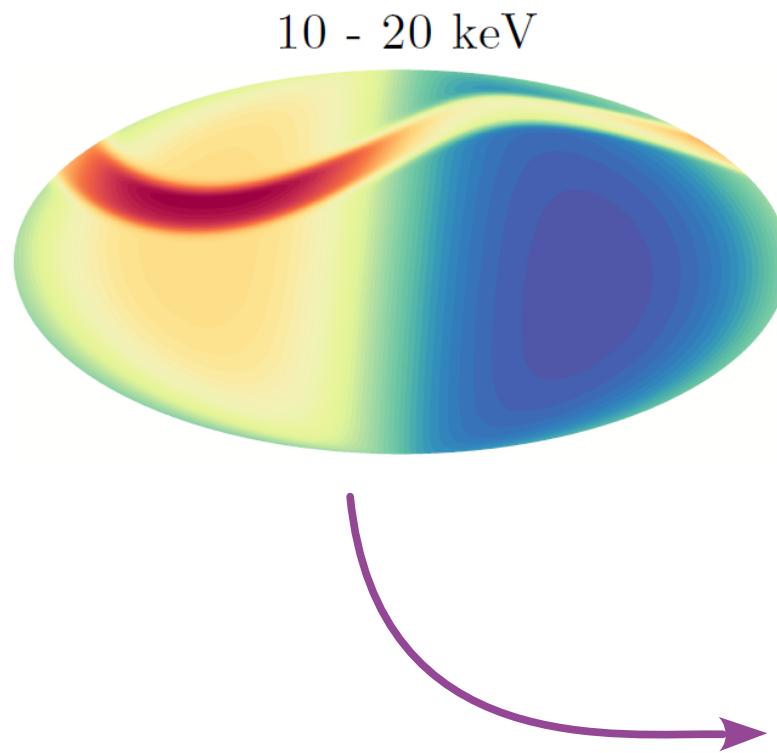
Detecting streams

- Tidal streams produce striking directional recoil patterns:

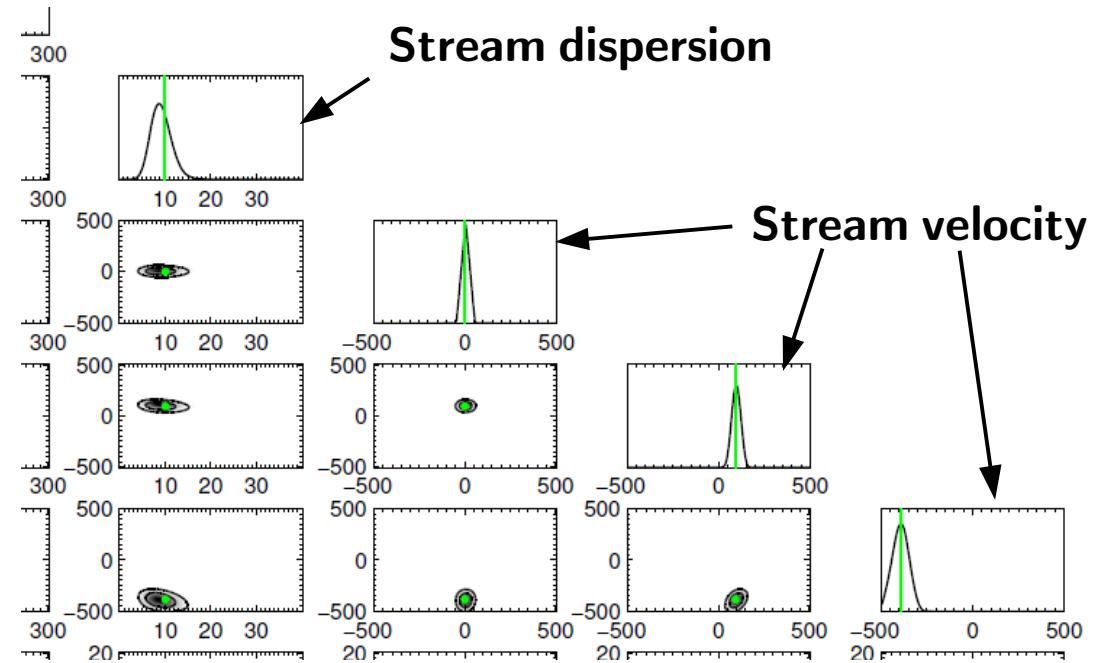


Detecting streams

- Could detect Sagittarius stream with 20 kg-year directional detector
 - Non-parametrically (test for median direction/rotational symmetry)
 - Parametrically (model stream → likelihood fit)



O'Hare & Green [1410.2749]



We want something more general...

There are ways to deal with the unknown speed distribution in standard *non-directional* detection e.g.,

- Halo independent methods
e.g. Fox [1011.1915], Frandsen [1111.0292], Kahlhoefer [1607.04418], and many more...
 - Measure $g(v_{\min}) = \int_{v_{\min}}^{\infty} \frac{1}{v} f(v) dv$ from data
- General parameterisations e.g. Peter [1103.5145], Kavanagh & Green [1303.6868]
 - use something like $f(v) = \exp\left(-\sum_{k=0}^{N-1} a_k v^k\right)$ and fit for a_k
- Add uncertainties on parameters into fit e.g. Strigari & Trotta [0906.5361]
 - pick a model and estimate astrophysical parameters
(with priors from measurements of MW)
- Add more parameters to deal with non-Maxwellian structure
e.g. Lee, & Peter [1202.5035], O'Hare & Green [1410.2749]

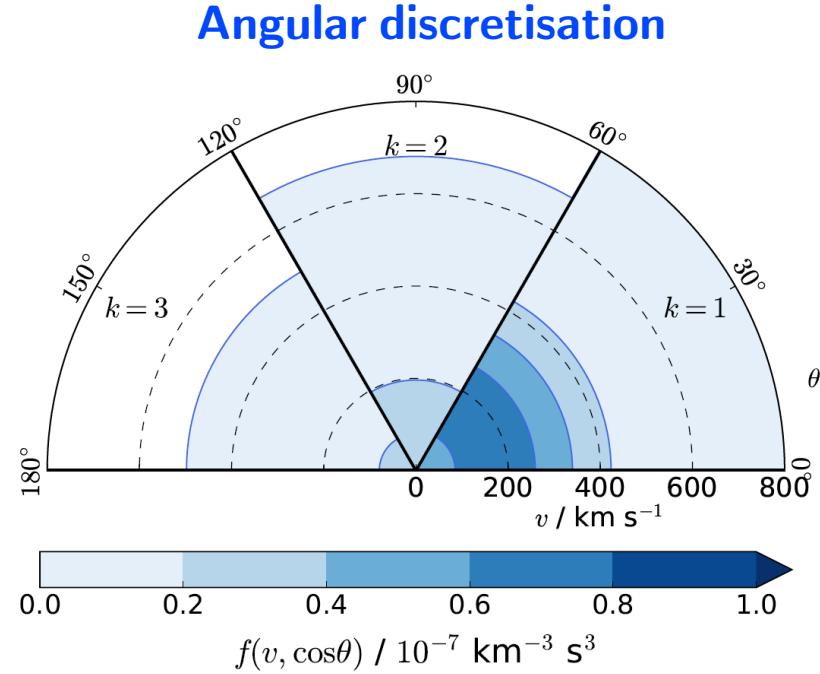
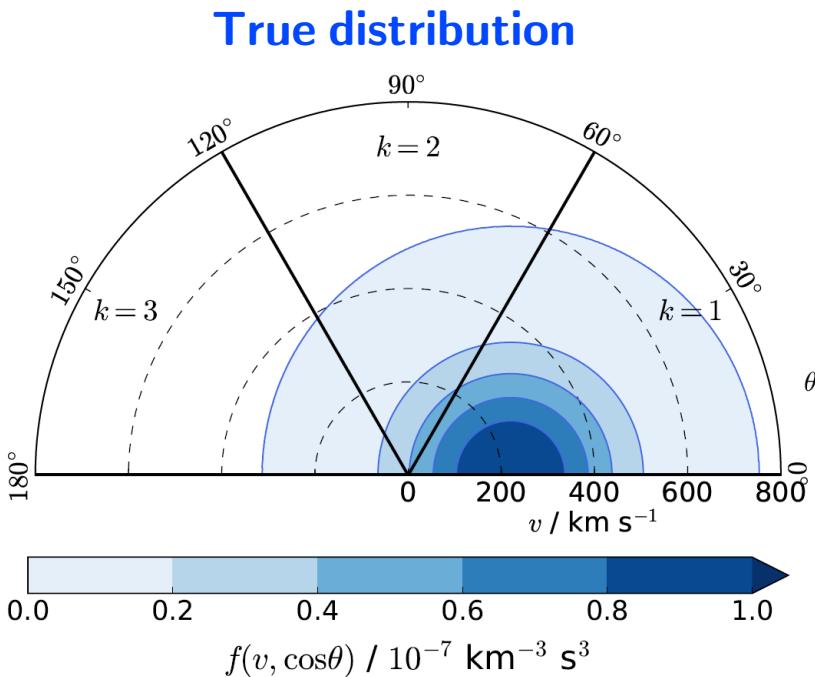
...But what about the velocity distribution?

Empirical velocity distribution

B. J. Kavanagh & C. A. J. O'Hare [1609.08630]

- Can we extract the velocity distribution from directional experiments in a model independent way?

$$f(\mathbf{v}) = f(v, \cos \theta, \phi) = \begin{cases} f^1(v) & \text{for } \theta \in [0^\circ, 60^\circ] \\ f^2(v) & \text{for } \theta \in [60^\circ, 120^\circ] \\ f^3(v) & \text{for } \theta \in [120^\circ, 180^\circ] \end{cases} \rightarrow \begin{array}{l} \text{Empirical polynomial} \\ \text{fit in each bin} \end{array} \rightarrow \text{Kavanagh [1502.04224]}$$



Reconstructing the velocity distribution

- For a given benchmark model generate mock data for two future directional detectors

	<u>Target</u>	<u>Threshold</u>	<u>Exposure</u>
<u>Experiment 1:</u>	Xe	5 keV	1 ton-year
+			
<u>Experiment 2:</u>	F	20 keV	10 kg-year

- Compare three methods of reconstruction

Method A:
Best case

We know the underlying velocity distribution and its parameters

Fit: mass, cross-section

Method B:
Reasonable case

We know the form of the velocity distribution but not the parameters

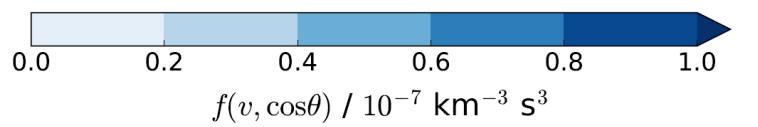
Fit: mass, cross section + astrophysical params.

Method C:
Worst case

We know nothing at all about the velocity distribution

Fit: mass, cross section + empirical parameters

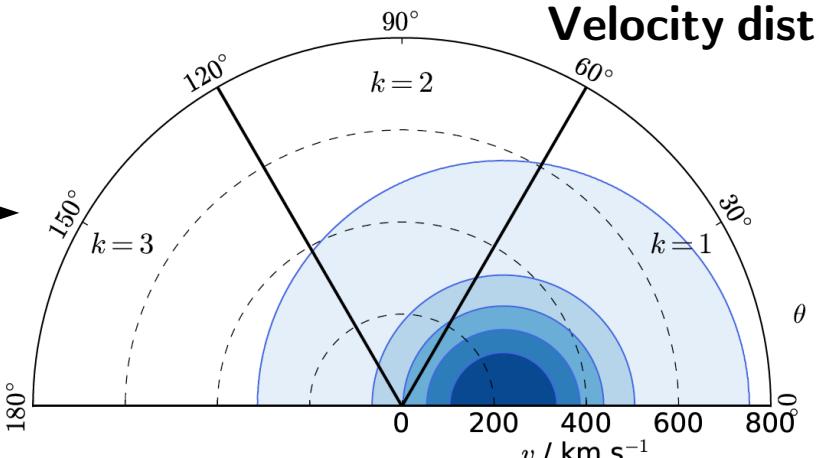
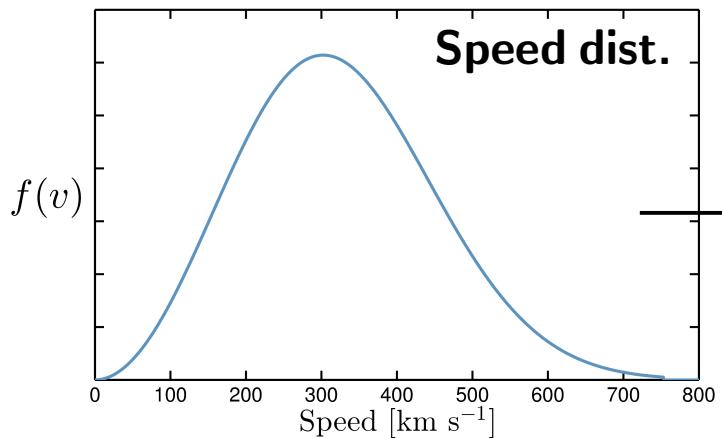
Benchmarks



SHM:

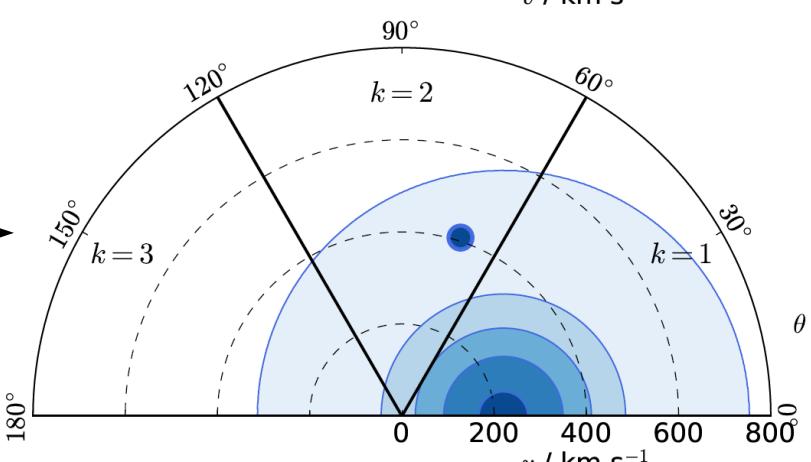
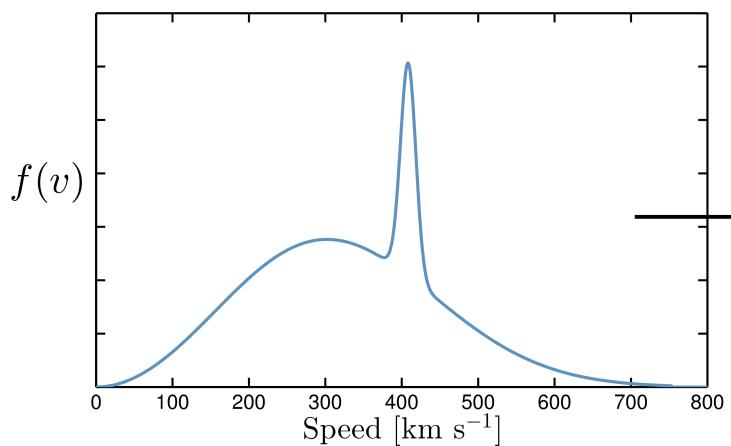
$$v_0 = 220 \text{ km/s}$$

$$v_{\text{esc}} = 533 \text{ km/s}$$



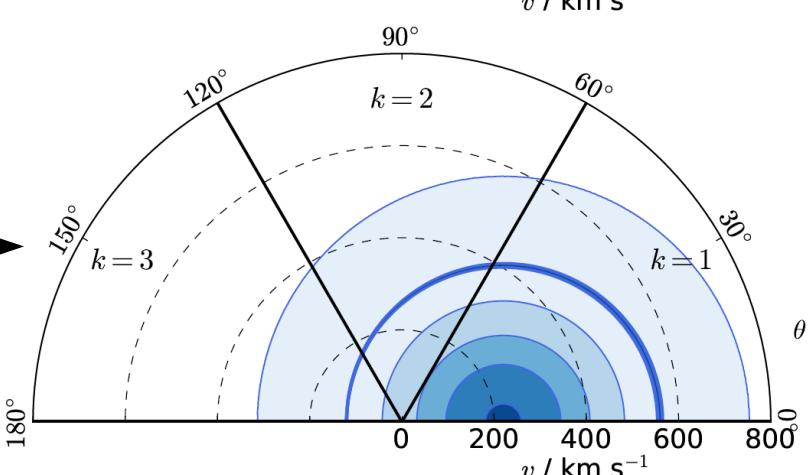
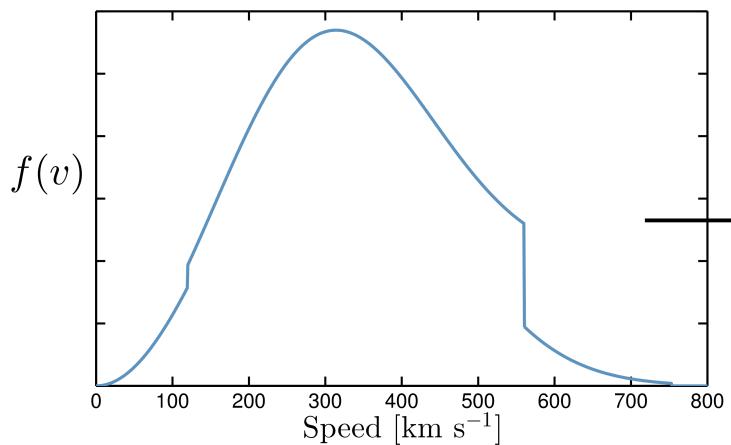
SHM+Stream:

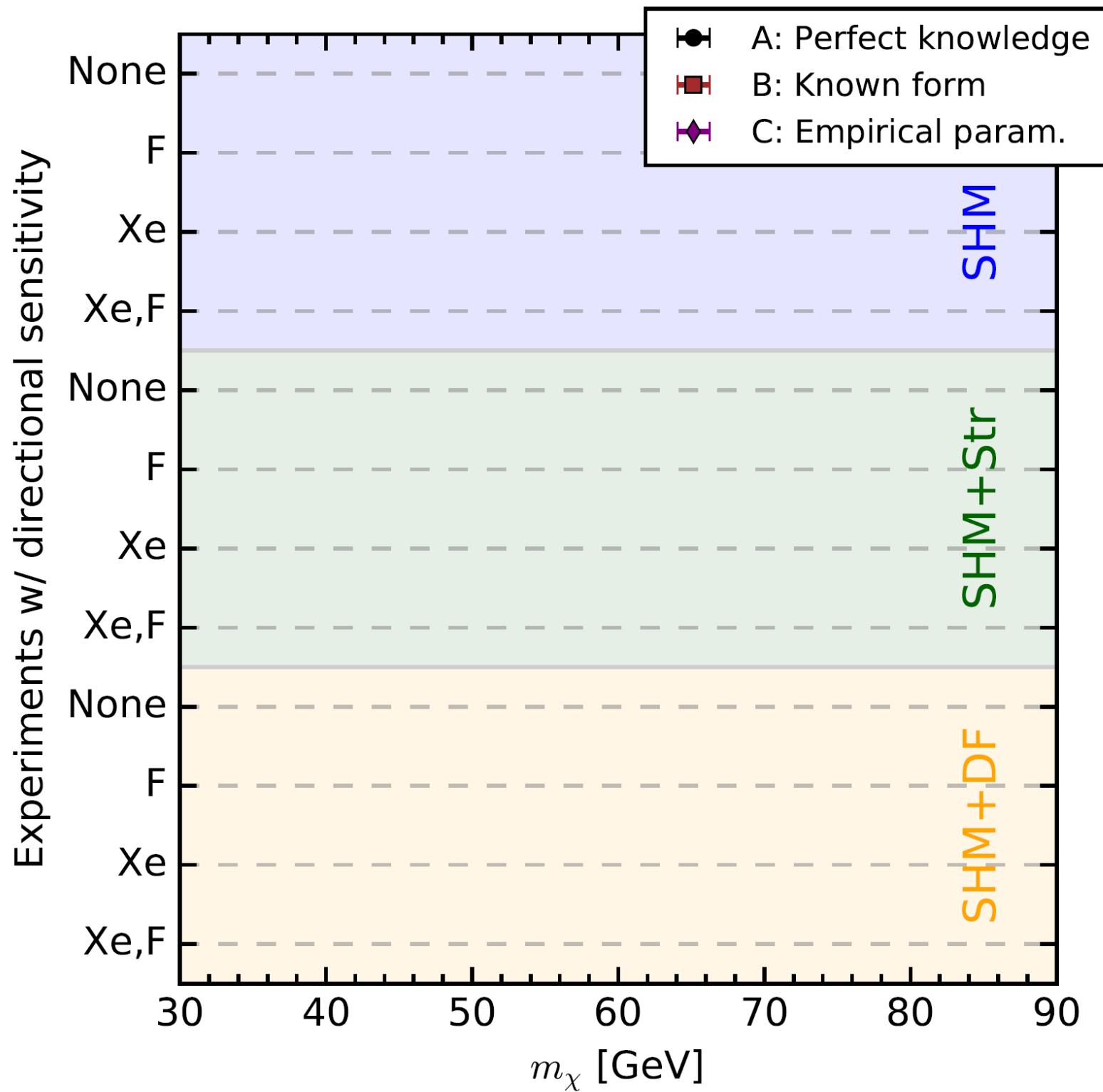
Purcell [1203.6617]

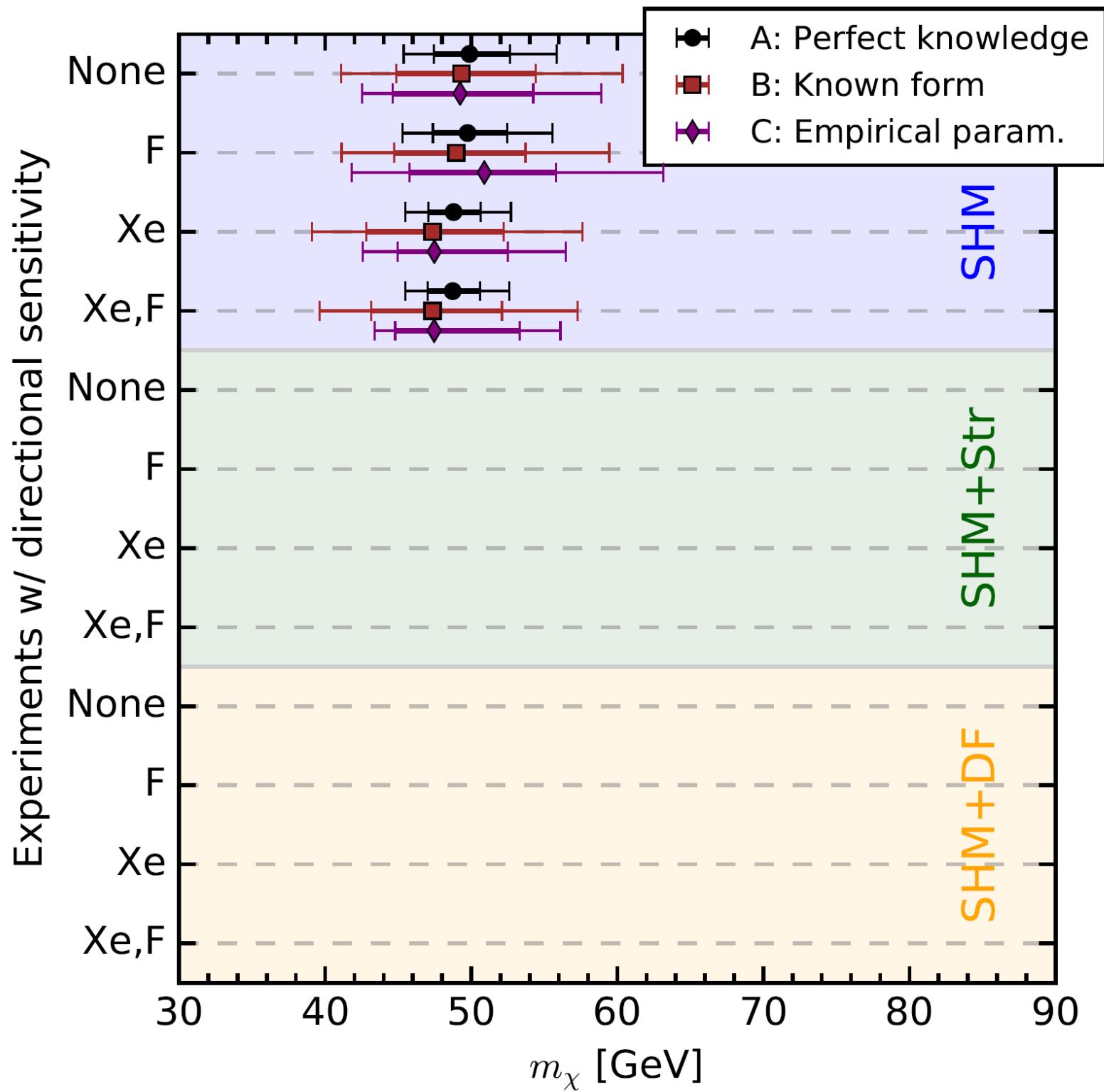


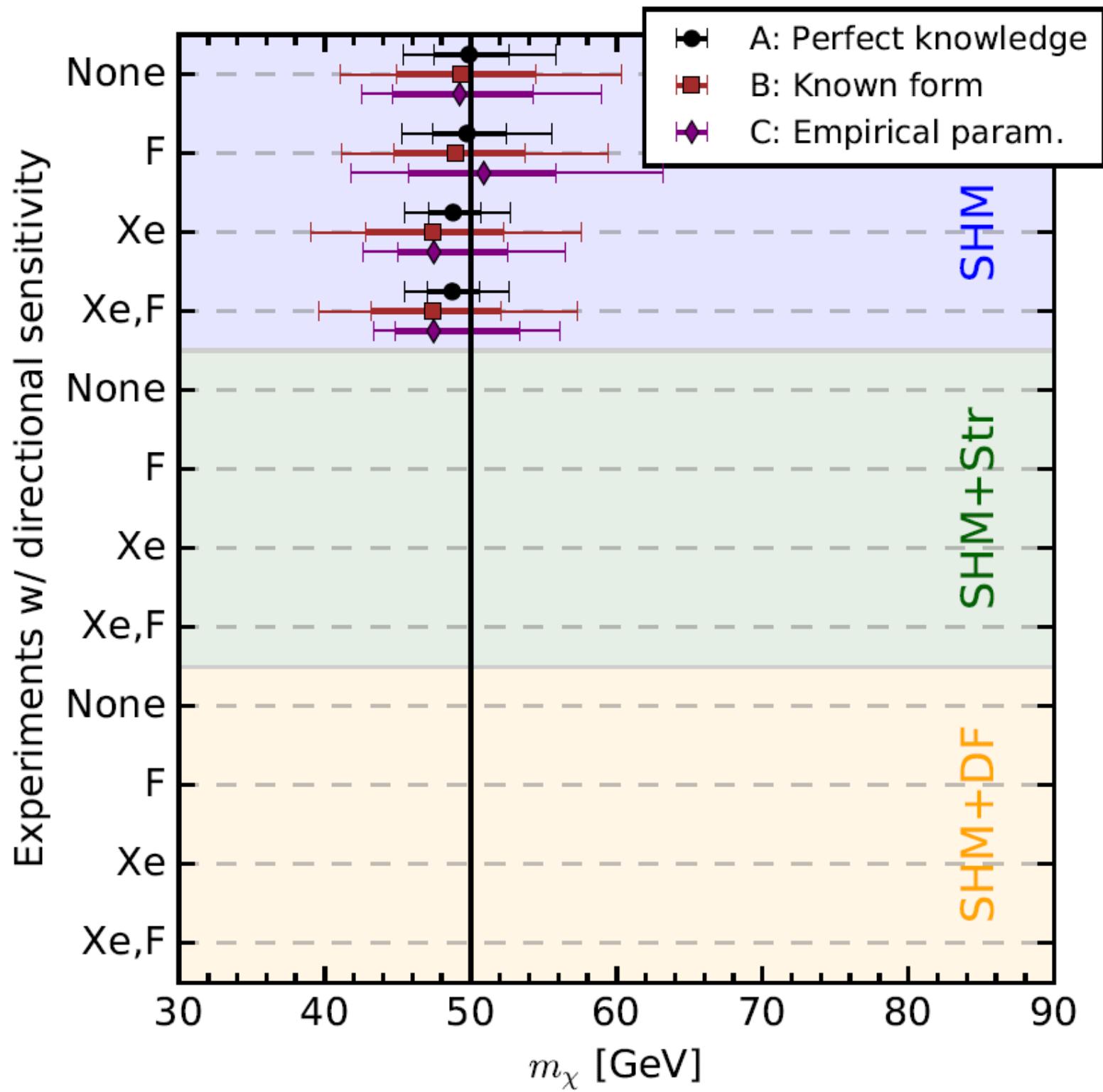
SHM+Debris flow:

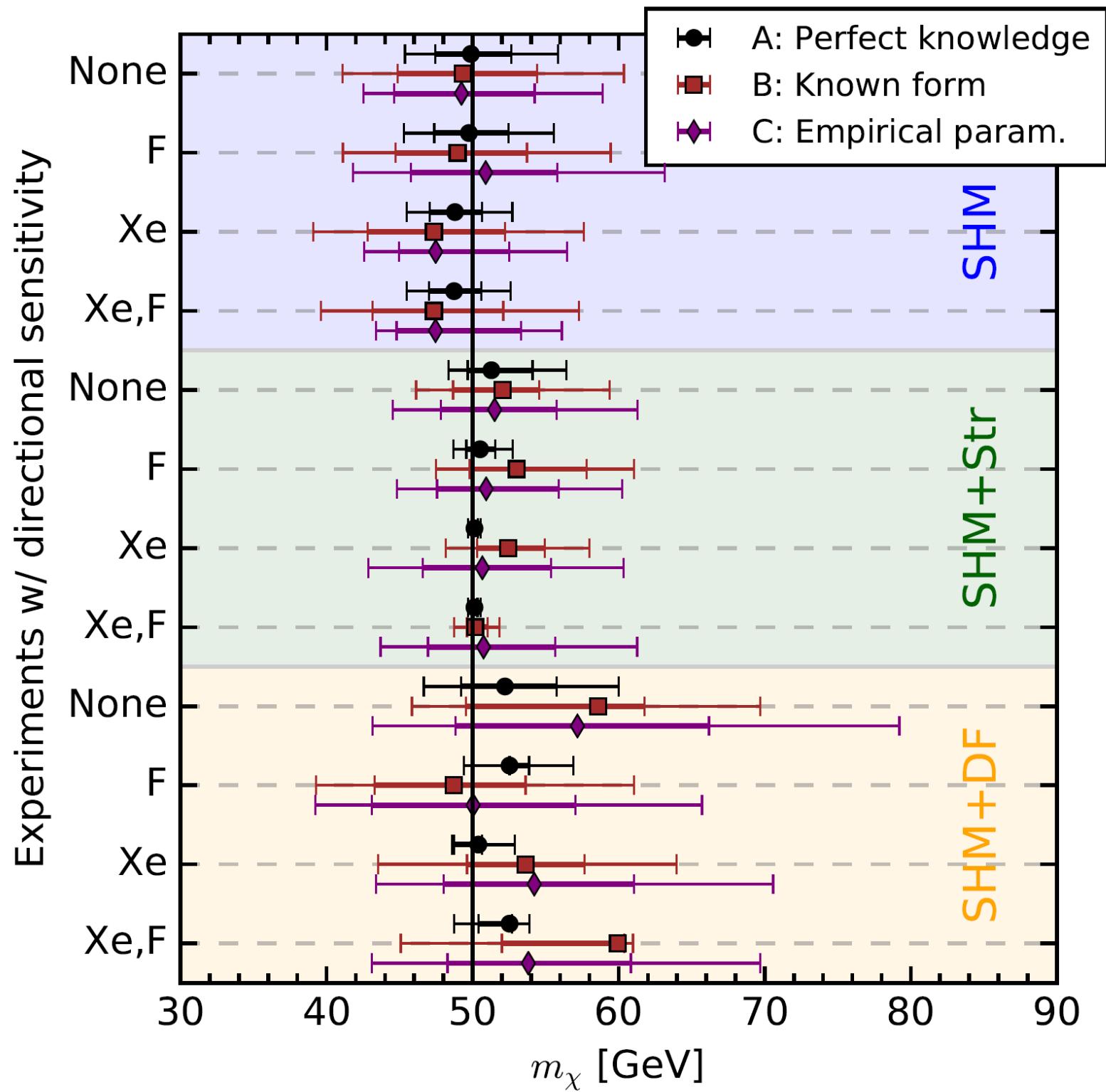
Kuhlen [1202.0007]



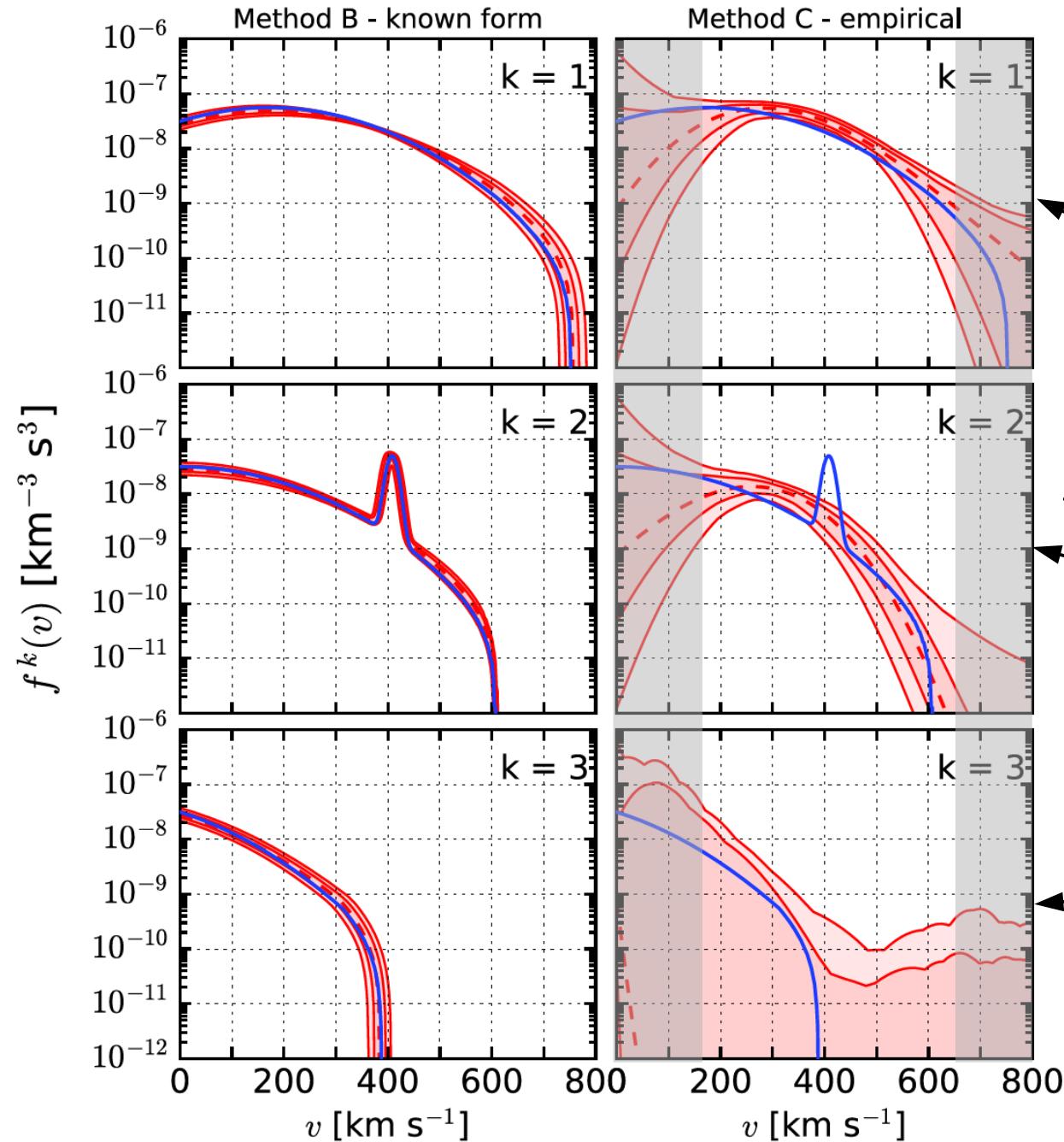








Shape of velocity distribution



Reconstructed binned distributions

→ SHM+Stream model

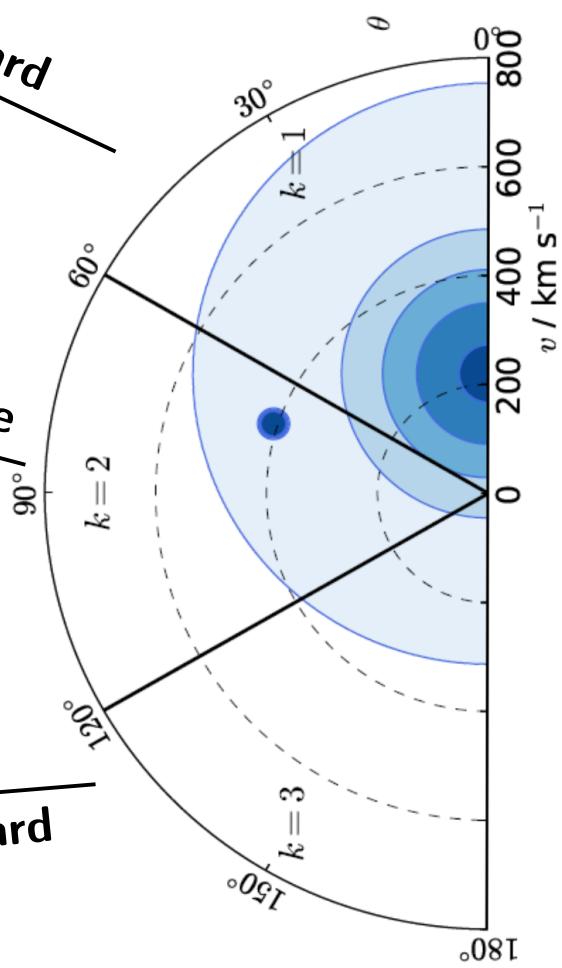
true distribution

best fit distribution

Forward

Transverse

Backward

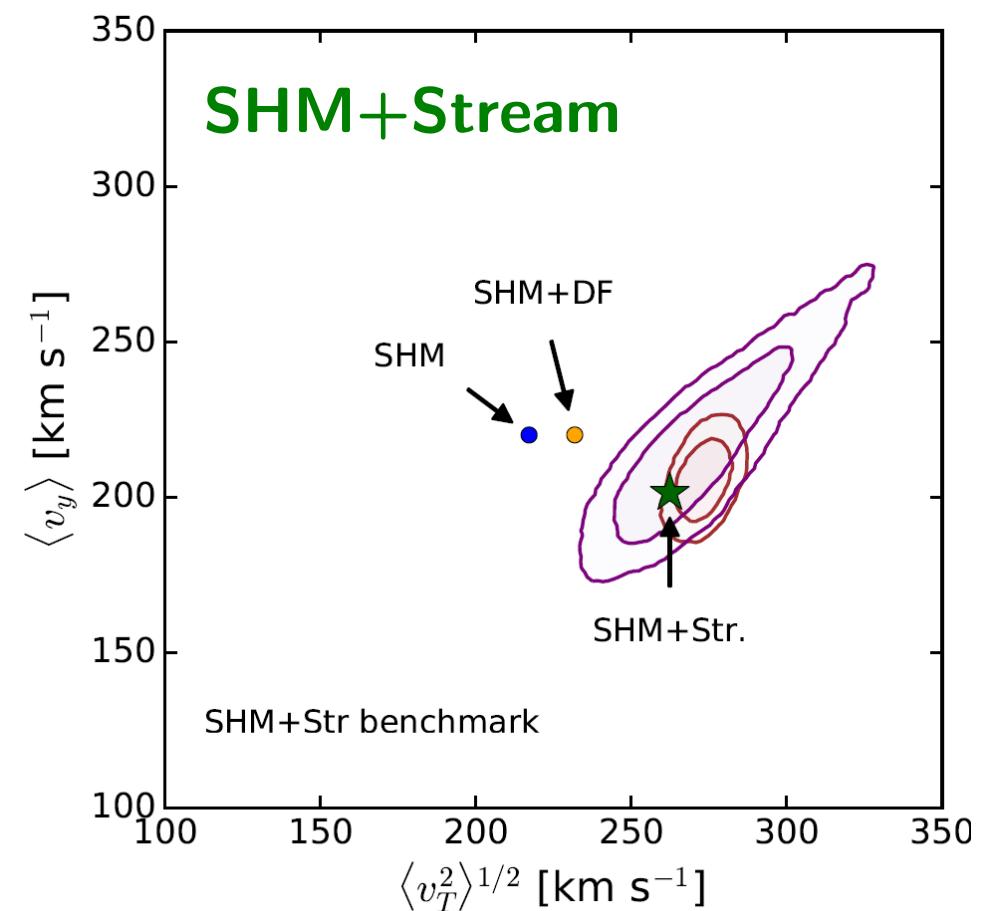
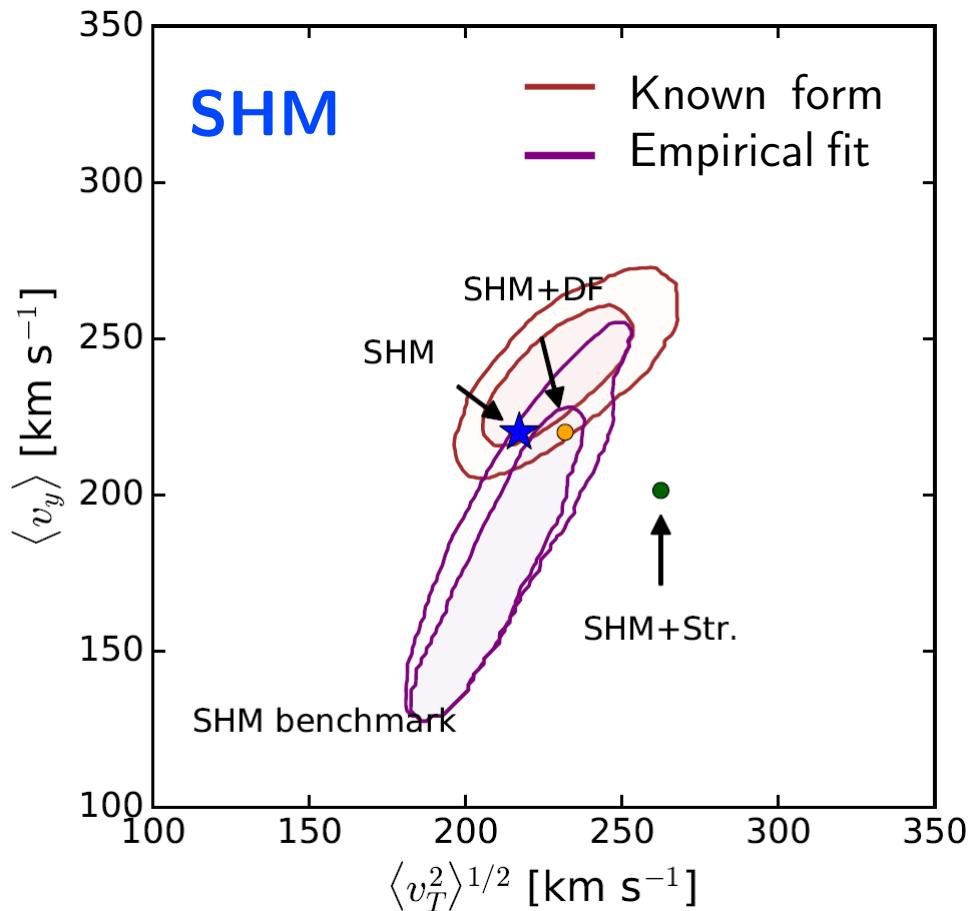


Comparing reconstructions

- Ideally we extract something physical to hint towards substructure
e.g. directionally averaged speeds:

$$\text{Average parallel to Earth's motion: } \langle v_y \rangle = \int dv \int_0^{2\pi} d\phi \int_{-1}^1 d \cos \theta (v \cos \theta) v^2 f(\mathbf{v})$$

$$\text{Average transverse to Earth's motion: } \langle v_T^2 \rangle = \int dv \int_0^{2\pi} d\phi \int_{-1}^1 d \cos \theta (v^2 \sin^2 \theta) v^2 f(\mathbf{v})$$



How could we use this method?

1. Detect dark matter (in a directional detector)
...
2. Measure WIMP parameters
→ use an empirical method to avoid bias
3. Derive velocity parameters
→ Do they hint towards substructure?
4. If so, fit with particular model
→ substructure detected?

Axions

Strong CP problem

QCD permits a CP violating term,

Gluon fields	Quarks	CP-violating term
$\mathcal{L}_{\text{QCD}} = -\frac{1}{4}G_{\mu\nu}^a G^{a\mu\nu} + \sum_q \bar{q}(i\gamma_\mu D^\mu - \mathcal{M}_q)q + \frac{\theta g^2}{32\pi^2} G_{\mu\nu}^a \tilde{G}^{a\mu\nu}$		

But neutron EDM
measurements \rightarrow
constrain phase

$$\bar{\theta} = \theta + \arg \det \mathcal{M}_q < \underline{\underline{10^{-10}}}$$

Baker *et al* [hep-ex/0602020]

Strong CP problem: why is the CP-violating
phase so unnaturally small?

Peccei-Quinn mechanism

- Solution to strong-CP problem: promote phase to dynamical field

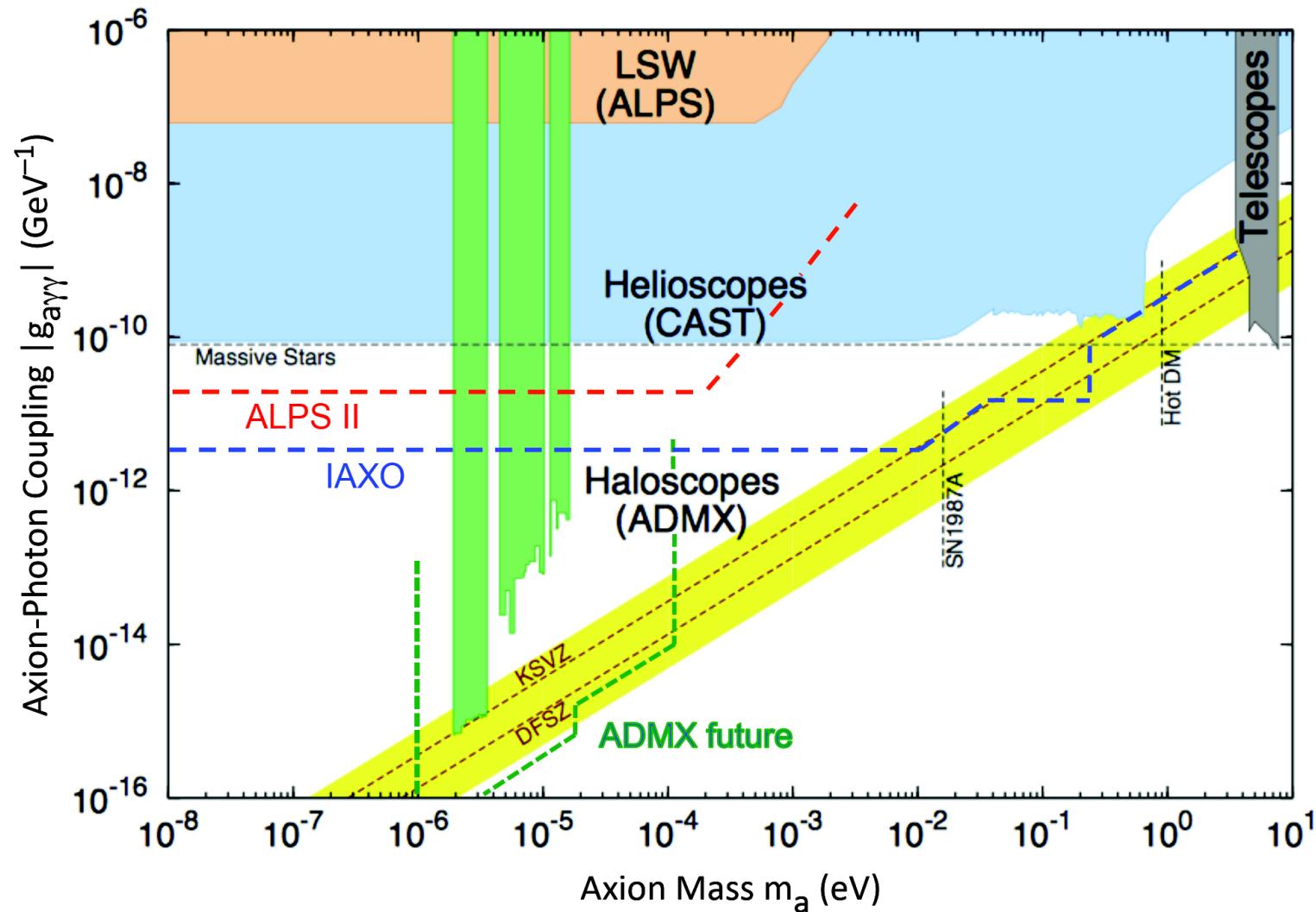
$$\mathcal{L}_{\text{QCD+axion}} = \dots + \frac{1}{2} \partial_\mu a \partial^\mu a + \frac{g^2}{32\pi^2} \frac{a(x)}{f_a} G_{\mu\nu}^a \tilde{G}^{a\mu\nu}$$

- Field vev nulls the CP violation
- Predicts a new particle, **the axion**
- Small mass given by QCD instanton effects (pion mixing)

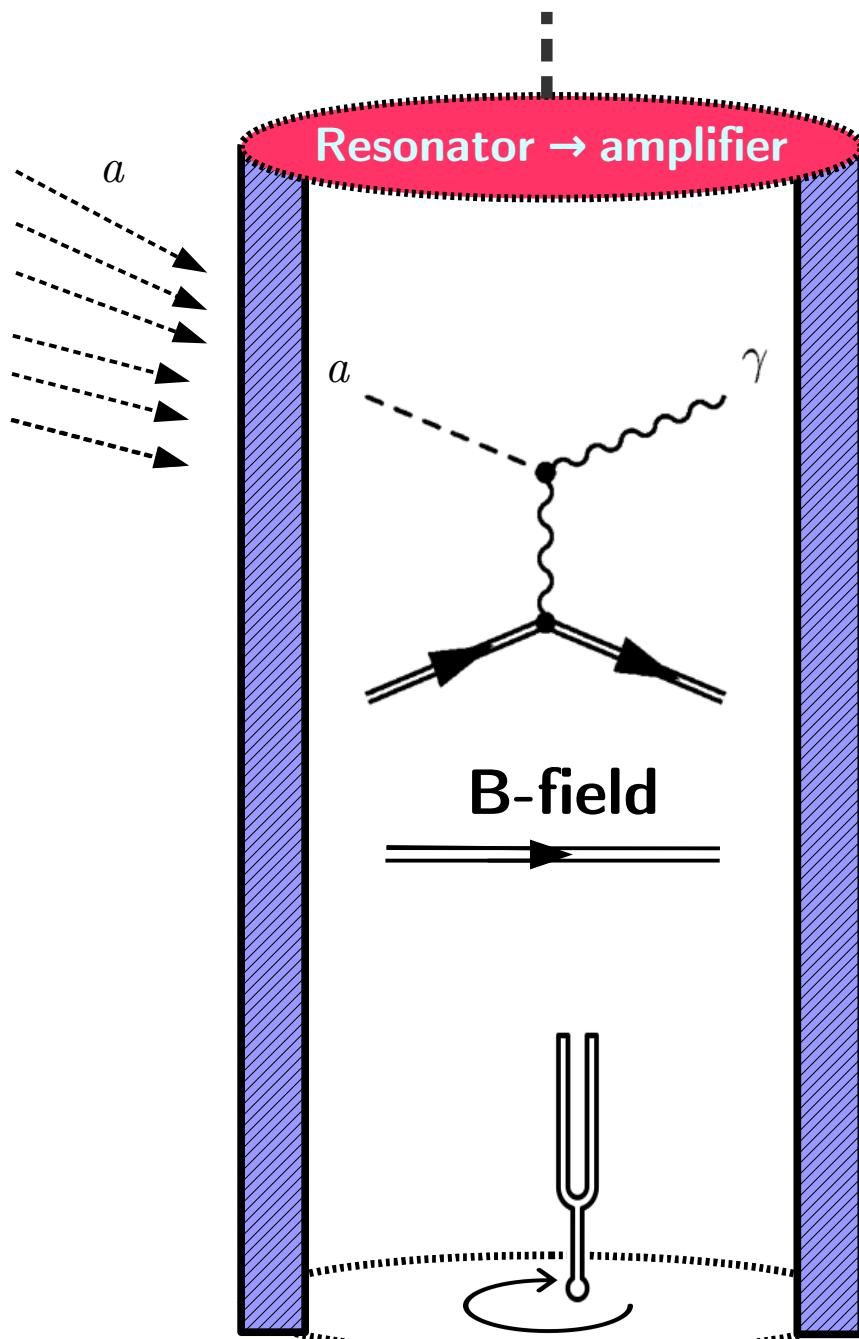
$$m_a \sim 6 \text{ eV} \left(\frac{10^6 \text{ GeV}}{f_a} \right)$$

Axion/ALP constraints

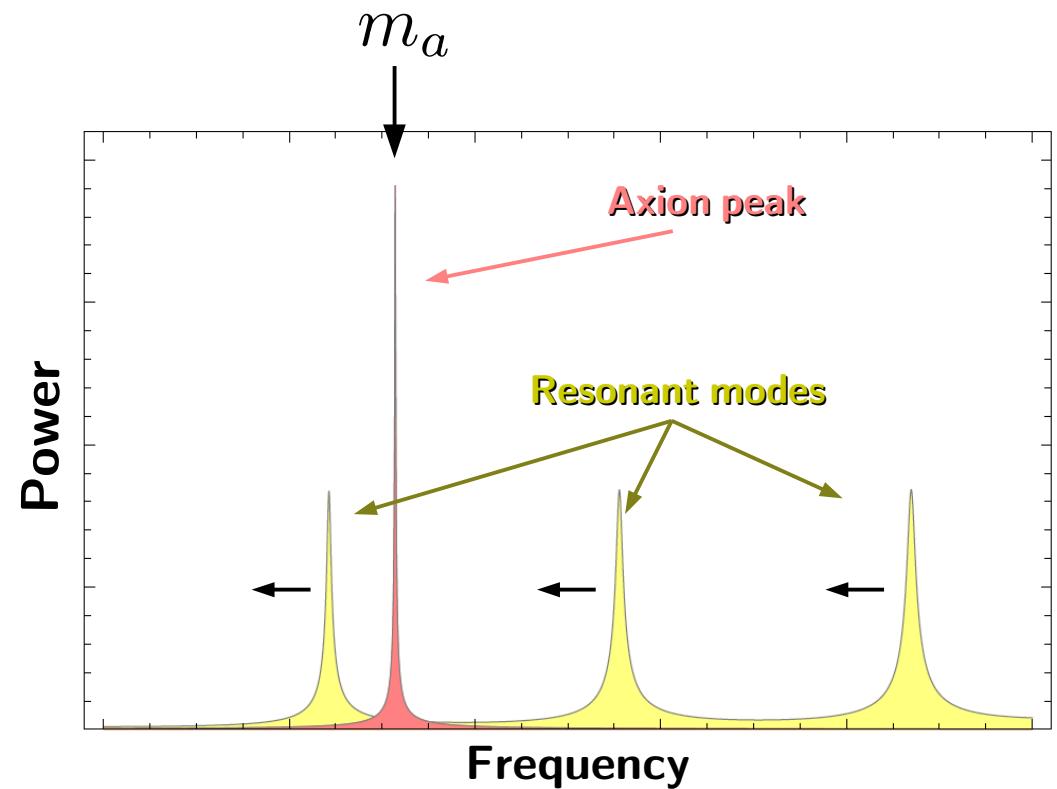
- Generalise to light particles outside of the PQ solution,
→ **axion-like particles (ALPs)**
- Measure conversion to photons inside magnetic fields:



Axion Haloscope



- **Signal enhancement:**
 - resonant frequency = axion mass
 - $$h\nu = m_a c^2$$
- **But we don't know axion mass...**
 - Must scan over a range of frequencies

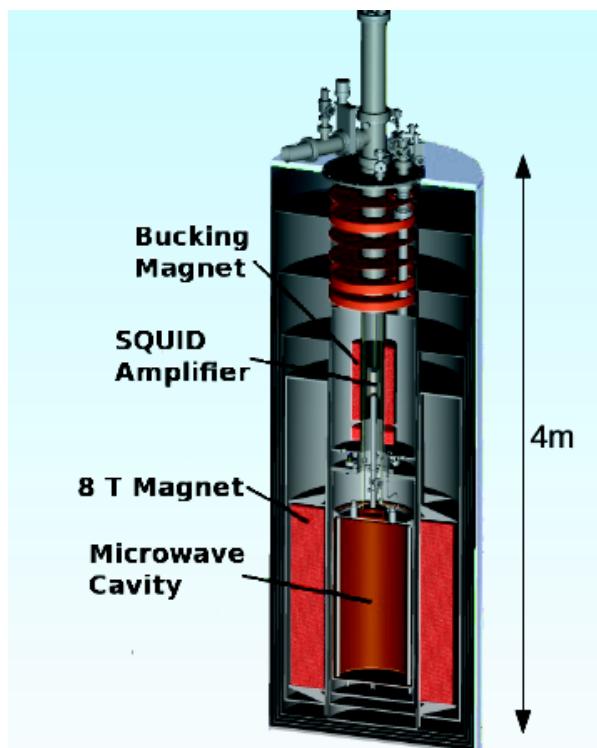


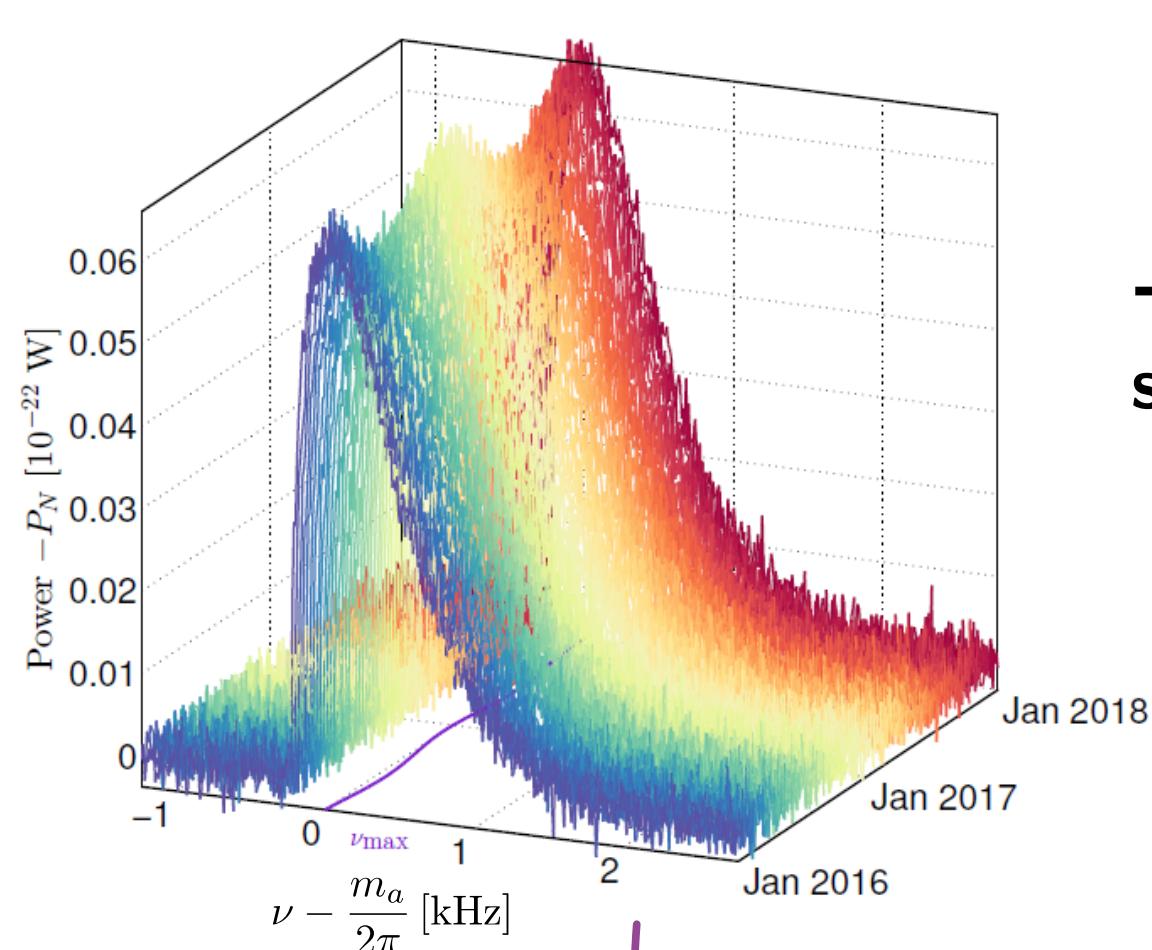
Axion Dark Matter Experiment (ADMX)



Asztalos et al [0910.5914]

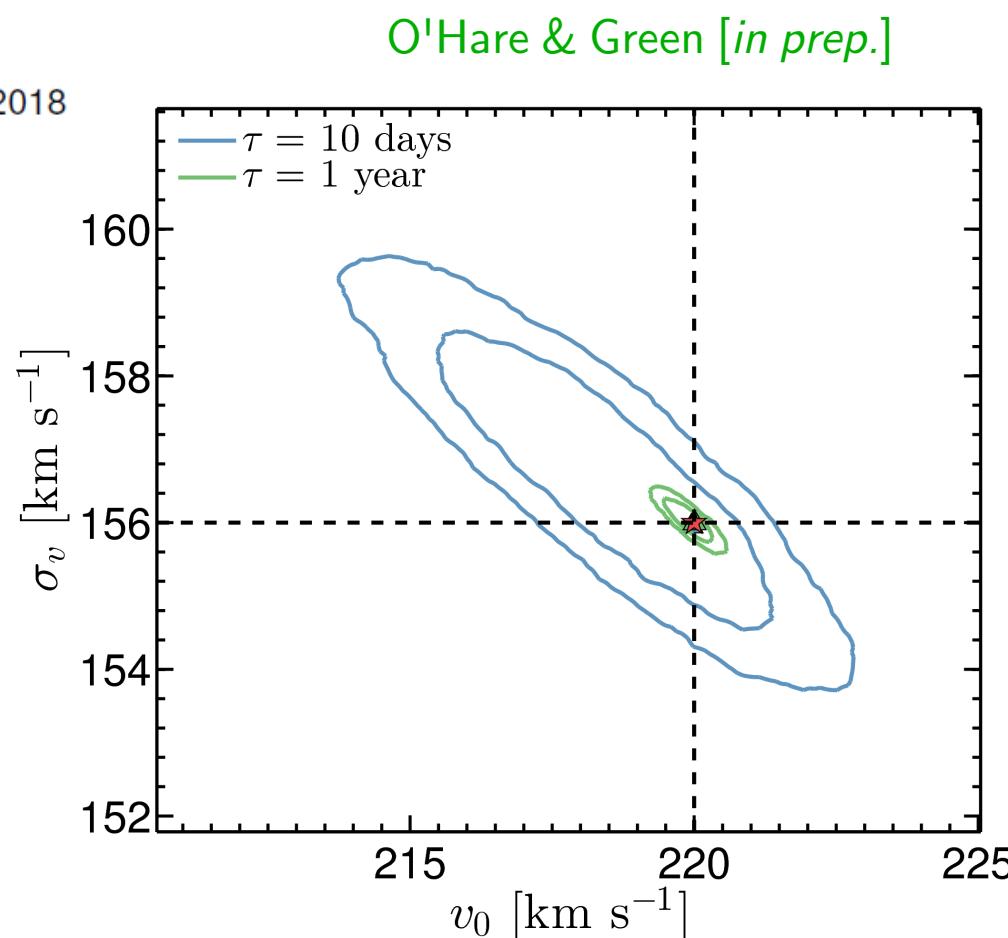
- Resonate at microwave freqs.
- 8T B-field
- 4 K noise temperature
- 200 litre volume
- Measure power $< 10^{-22} \text{ W}$



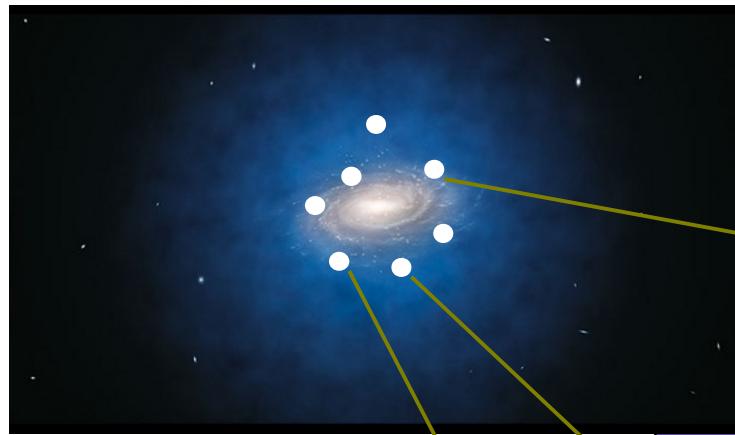


Extract
astrophysical
parameters

→ **Measure axion power spectrum**

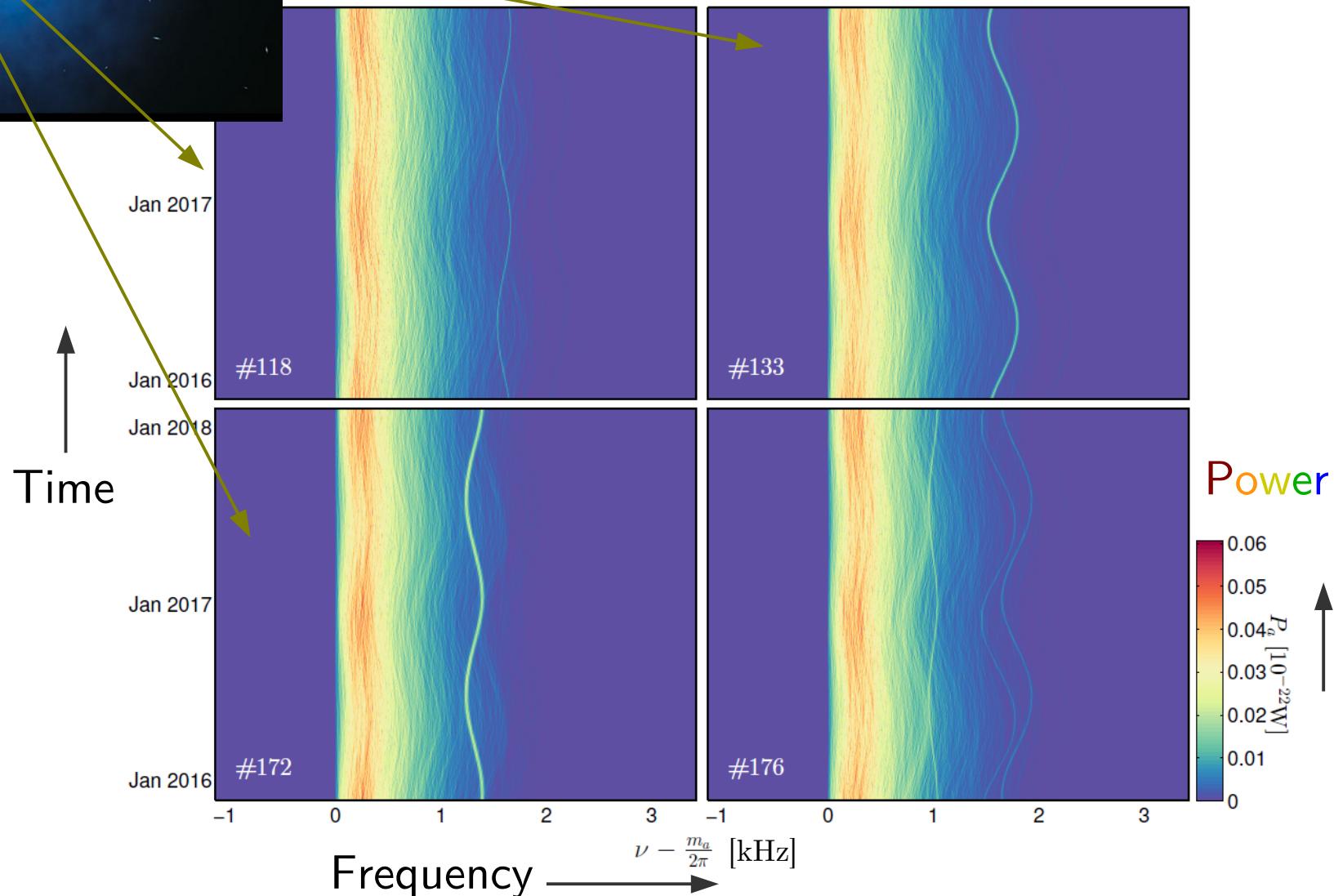


Example power spectra from distributions in VL2 simulation



- Samples from 1 kpc bubbles @ 8 kpc Gal. radius

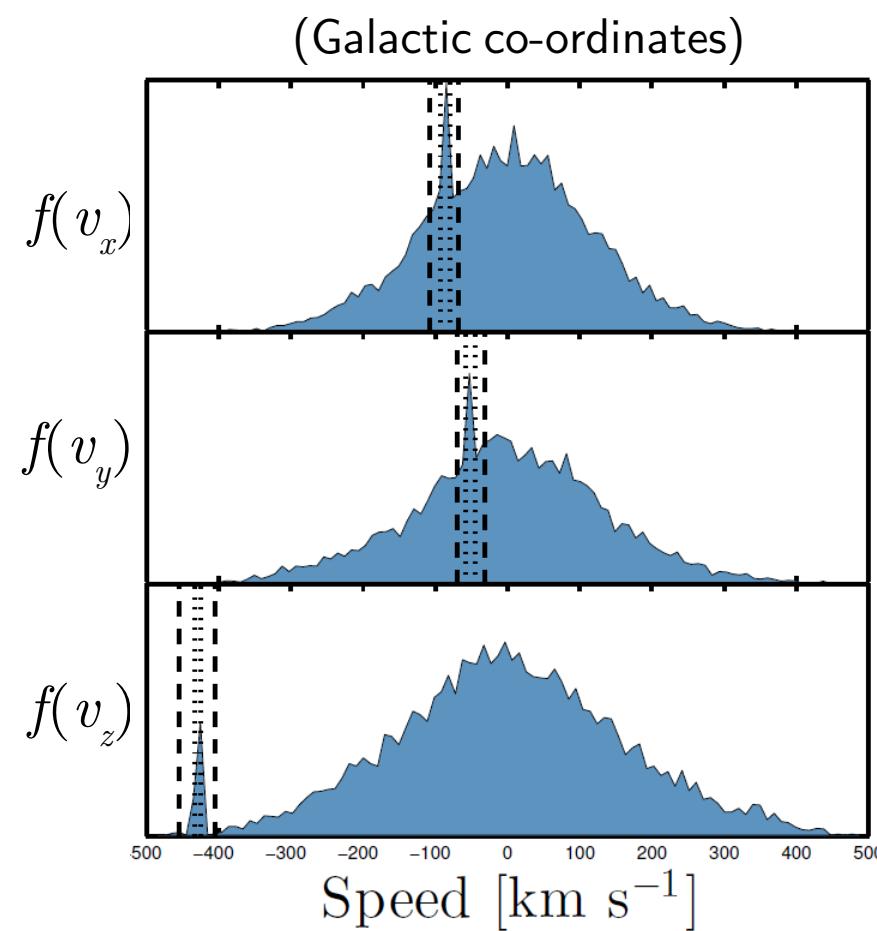
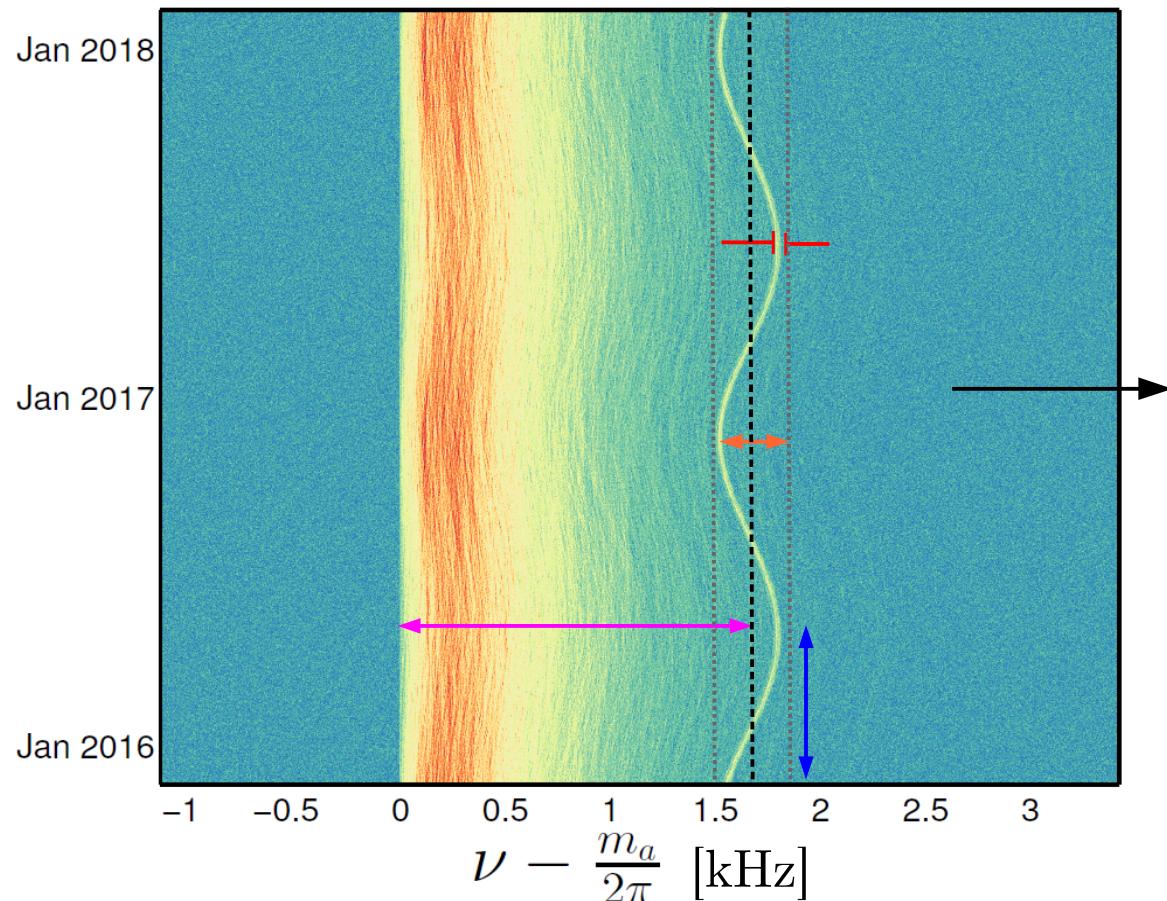
O'Hare & Green [*in prep.*]



Measuring axion streams

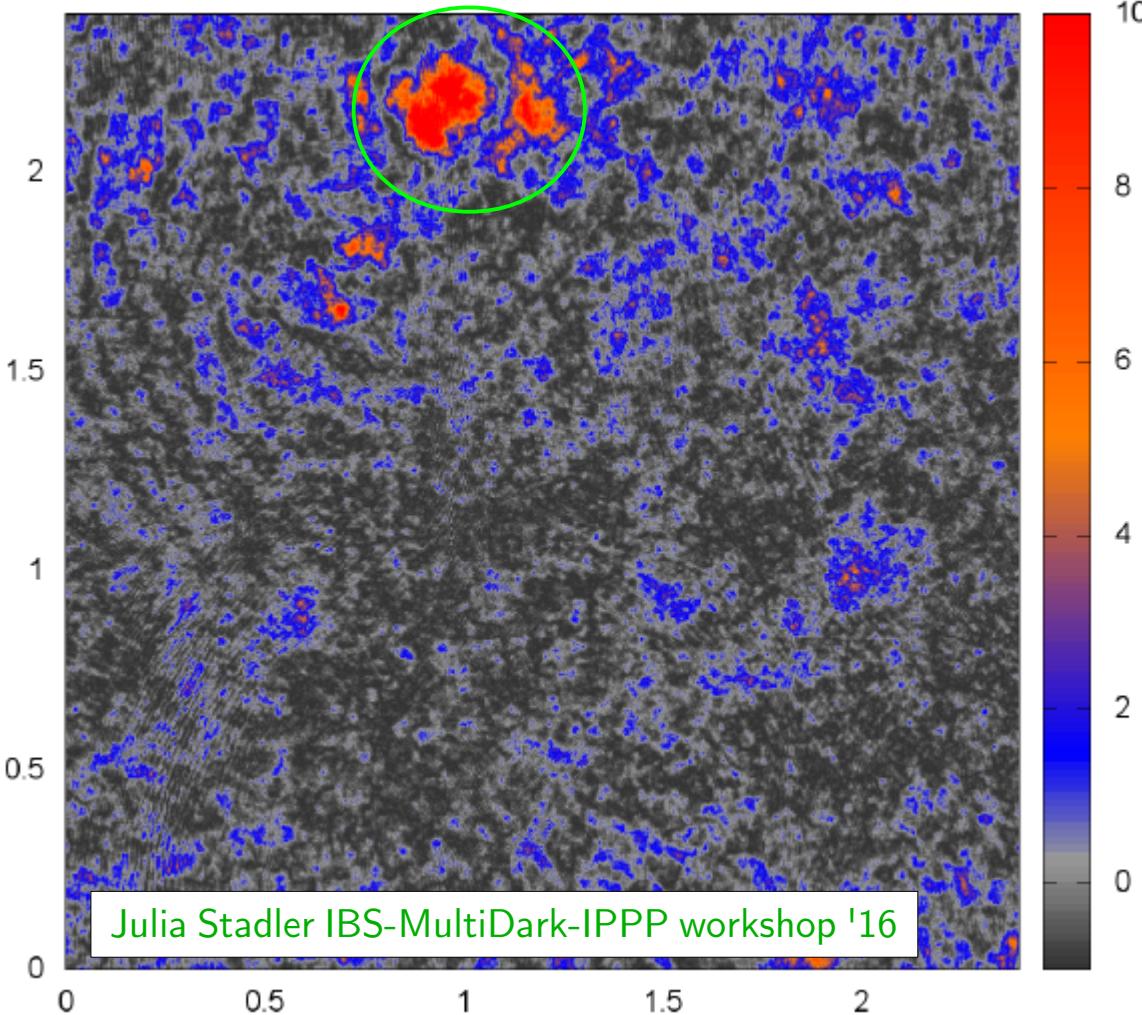
- Can extract all five properties of a stream from sinusoid
 - Stream density ← **Power** relative to bulk
 - Stream dispersion ← **Width** of sinusoid
 - Galactic velocity ← **Amplitude, phase and mean** of sinusoid

O'Hare & Green [*in prep.*]



Axion miniclusters

- Collapsing density perturbations in the early Universe can form small clusters of axions



Density contrast: $\Phi = \frac{\delta\rho}{\rho} \sim 1$

Mass: $M \sim 10^{-12} M_\odot$

Radius: $R \sim 10^7$ km

Density: $\rho \sim 10^6$ GeV cm $^{-3}$

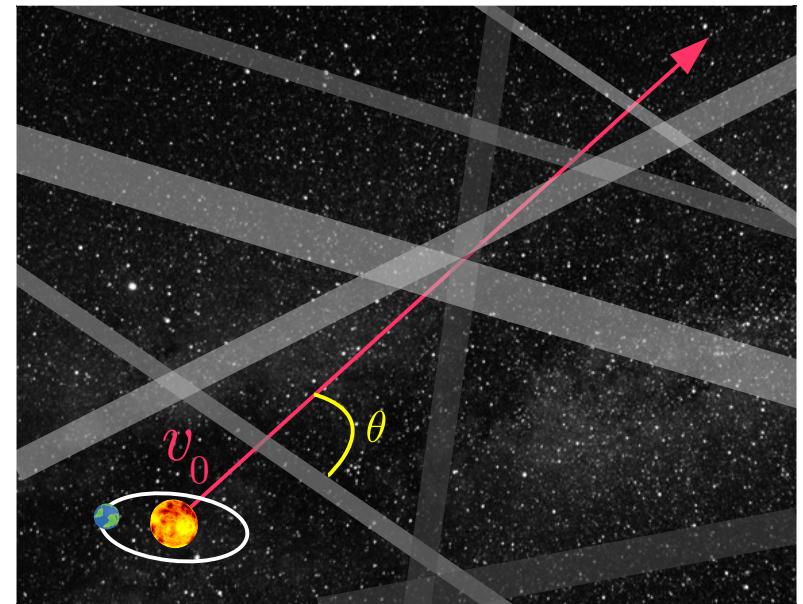
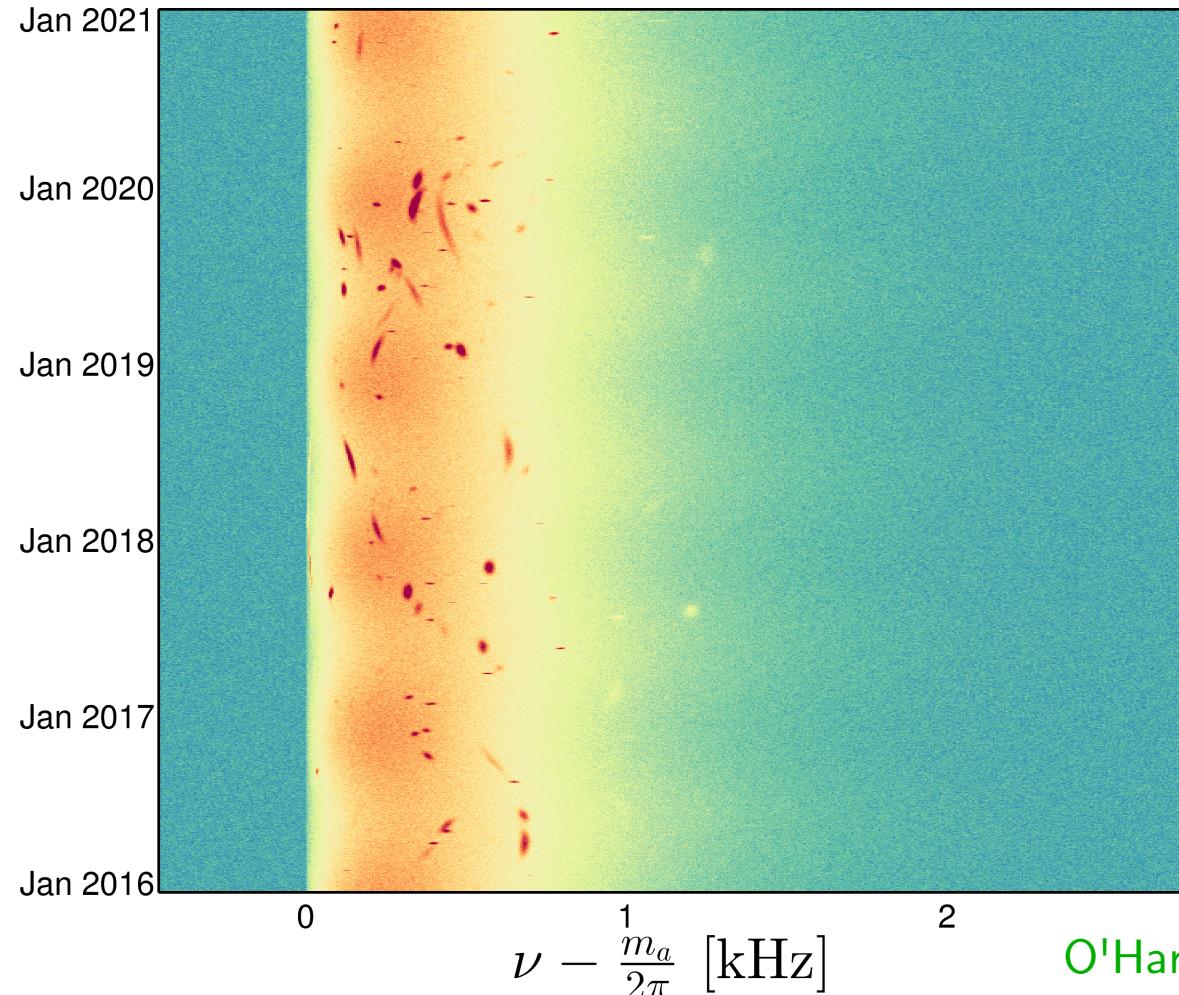
- Could comprise non-negligible fraction of DM halo
- Up to 10^{10} pc $^{-3}$ in the local stellar neighbourhood?

Tinyakov [1512.02884]

Axion miniclusters

- Tidally disrupted by interactions with stars
 - network of minicluster streams wrapping Milky Way
 - stream crossing time $\mathcal{O}(1\text{-}100)$ days Tinyakov [1512.02884]

$$R_{\text{str}} \simeq \frac{3.4 \times 10^7 \text{ km}}{\Phi(1 + \Phi)^{1/3}} \left(\frac{M_{\text{mc}}}{10^{-12} M_{\odot}} \right)^{1/3} \longrightarrow \tau_{\text{str-x}} = \frac{2R_{\text{str}}}{v_0 \sqrt{1 - \frac{\mathbf{v}_{\text{str}} \cdot \mathbf{v}_{\text{lab}}}{v_0 v_{\text{str}}}}}$$



Would need to be separated from environmental noise peaks with use of time/daily modulation

Summary

**We must understand
local MW halo to do dark
matter detection**

- Uncertainty in exclusion limits
- Biased particle measurements
- Neutrino floor closer to existing limits

**Dark matter detectors
can help us understand
the local MW halo**

- Directional detectors/
haloscopes well suited to
study the local halo
- Observe non-Maxwellian
structure
- Learn about formation
history of MW

Dark matter halo

Event rate dependent on $\rho_0 \int f_{\text{lab}}(\mathbf{v}, t) \delta(\mathbf{v} \cdot \mathbf{q} - \hat{v}_{\min}) d^3 \mathbf{v}$

- local density ρ_0
- velocity distribution $f_{\text{lab}}(\mathbf{v}, t) = f_{\text{gal}}(\mathbf{v} + \mathbf{v}_{\text{lab}}(t))$

→ lab velocity: $\mathbf{v}_{\text{lab}}(t) = \mathbf{v}_{\text{LSR}} + \mathbf{v}_{\text{pec}} + \mathbf{v}_{\text{rev}}(t) + \mathbf{v}_{\text{rot}}(t)$
 $\sim 220 \text{ km/s}$ $\sim 18 \text{ km/s}$ $\sim 30 \text{ km/s}$ $\sim 0.5 \text{ km/s}$

